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# THE BEHAVIOR OF THREE NODE POWER NETWORKS 

by

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The Behavior of Three Node Power Networks

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#### Abstract

The real power flow equation for a three node network is analyzed in terms of its topological and geometric aspects. It is shown that the set of feasible power injections is convex and to each feasible injection there corresponds a unique stable solution. Certain aspects of the behavior of the associated swing equations are studied. This behavior is quite different from that of a two node network; in particular, the system may not be completely stable even when the damping is made arbitrarily large.

\section*{1. Fntroduction}

While great advances have been achieved in the numerical solution of the power flow equation, important questions pertaining to the geometric and topological character of the solutions remain unresolved (see [1].).

This study presents a comprehensive analysis of three node power networks in which the transmission network is modeled by the real power flow equation and the generator dynamics by the classical swing equations. Primary attention is devoted towards revealing the geometric and topological structure of the power flow equation.

It is clear from the literature that intuition guiding both load flow and transient stability studies derives largely from the well-understood one dimensional problem, i.e., a two node network. It turns out, however, that even in the "slightly" more complex case of three nodes, the behavior changes dramatically and conjectures that seem reasonable in light of the one dimensional problem are invalid.

Analysis of the power flow equation for three nodes is quite involved and points to the great diffculties to be encountered in the general case. To the extent possible, however, we have used general arguments avoiding the topological properties of the plane.


Section 2 is devoted to the load flow equation. There are two principal results. First, the set of feasible power injectons is convex (Theorem 2.5). Second, to each feasible injection there corresponds a unique stable solution (Theorem 2.4). Surprisingly, the set of stable solutions need not be convex (Theorem 2.3).

Section 3 discusses certain properties of the associated pair of coupled swing equations. The main result is an example where the system is not completely stable no matter how large the damping is. This is in sharp contrast with the behavior of an isolated swing equation. The example is "generic" in the sense that every three node network loses complete stability whenever the power injections become sufficiently large.

## 2. Global Proprties of the Power Flow Equation

The principal properties of the power flow equation are derived in this section.

### 2.1 The Load Flow Function

Consider the three node network of Figure 1. Each node is a PV bus, and the third bus is a slack bus. Take the third bus voltage angle as reference and denote the voltage phasor at bus i by:

$$
\mathrm{V}_{1} \exp \left(\mathrm{j} v_{1}\right), \mathrm{i}=0,1,2 .
$$

By definition,

$$
v_{0} \equiv 0 .
$$

$Y_{i j}$ is the admittance of the lossless transmission line joining $i$ and $j$. Defining $B_{i j}=V_{i} V_{j} Y_{i j}$ we obtain the load flow equation in the form:

$$
\begin{align*}
& \mathrm{P}_{1}=\mathrm{f}_{1}\left(v_{1}, v_{2}\right)=\mathrm{B}_{1} \sin \left(v_{1}\right)+\mathrm{B}_{12} \sin \left(v_{1}-v_{2}\right),  \tag{2.1}\\
& \mathrm{P}_{2}=\mathrm{f}_{2}\left(v_{1}, v_{2}\right)=\mathrm{B}_{2} \sin \left(v_{2}\right)+\mathrm{B}_{12} \sin \left(v_{2}-v_{1}\right) .
\end{align*}
$$

Throughout this section we assume $B_{1}>0, B_{2}>0, B_{12}>0$; otherwise, the analysis is trivial in the sense that it reduces to the one dimensional case. Define the vector $v=\left(v_{1}, \vartheta_{2}\right)$. Utilizing the periodicity of function $f$, we restrict its domain to the set:

$$
\mathrm{T}^{2}=[-\pi, \pi]^{2}
$$

where the end points $\pi$ and $-\pi$ are identified. The Jacobian of f at $\vartheta$ is given by

$$
\mathrm{F}(v)=\frac{\partial \mathrm{f}}{\partial v}=\left[\begin{array}{cc}
\mathrm{B}_{1} \cos v_{1}+\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right) & -\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right)  \tag{2.2}\\
-\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right) & \mathrm{B}_{2} \cos v_{2}+\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right)
\end{array}\right]
$$

Following [1] write $F(v)>0$ or $F(v) \geq 0$ according as $F(v)$ is positive definite or positive semidefinite. Furthermore, $v \in \mathrm{~T}^{2}$ is said to be stable if $\mathrm{F}(\vartheta)>0$. Denote the stable region by $\mathrm{H}_{\mathrm{s}}$. In addition let $\mathrm{H}_{0}$ denote the subset of $v$ in $\mathrm{T}^{2}$ for which $F(v) \geq 0$ but $F(v)$ is not strictly positive definite.
Definition 2.1 (see [1]) The principal polytope $H_{p}$ is the subset of all $v \in \mathrm{~T}^{2}$ such that $\left|v_{1}\right| \leq \frac{\pi}{2},\left|v_{2}\right| \leq \frac{\pi}{2}$ and $v_{1}-v_{2} \left\lvert\, \leq \frac{\pi}{2}\right.$.

Also define the following subsets of $\mathrm{T}^{2}$ :

$$
\begin{align*}
& A_{1}=\left\{v \in T^{2}\left|v_{1}\right| \leq \frac{\pi}{2}, v_{1}-v_{2} \left\lvert\, \leq \frac{\pi}{2}\right.\right\}  \tag{2.3a}\\
& A_{2}=\left\{v \in T^{2}| | v_{2}\left|\leq \frac{\pi}{2}, v_{1}-v_{2}\right| \leq \frac{\pi}{2}\right\}  \tag{2.3b}\\
& A_{3}=\left\{v \in T^{2}\left|v_{1}\right| \leq \frac{\pi}{2}, \hat{w}_{2} \leq \frac{\pi}{2}\right\} \tag{2.3c}
\end{align*}
$$

Denote by $\mathrm{H}_{\mathrm{b}}$ the union of these subsets:

$$
\begin{equation*}
H_{b}=A_{1} \cup A_{2} \cup A_{3} \tag{2.4}
\end{equation*}
$$

These sets are exhibited in Figure 2.

### 2.2 The Geometry of $\mathrm{H}_{\mathrm{B}}$

We begin by establishing an estimate of the stable region.

Theorem 2.1 $\mathrm{H}_{\mathrm{p}} \subset \mathrm{H}_{\mathrm{s}} \cup \mathrm{H}_{0} \subset \mathrm{H}_{\mathrm{b}}$.
Proof: The first inclusion is proved in [1] for a general network. To prove the secpnd inclusion suppose that $v$ is not in $H_{b}$. Then of course $v$ is not in $A_{g}$. Hence $\cos v_{1}<0, \cos v_{2}<0$. Let $u=(1,1)$. Then $\langle u, F(v) u\rangle=B_{1} \cos v_{1}+$ $\mathrm{B}_{2} \cos \vartheta_{2}<0$ implying that $F(\vartheta)$ is not positive semidefinite. The assertion follows.

## Corollary 2.1

Let $\vartheta \in H_{s} \cup H_{0}$. Then for $0 \leq \varepsilon<1, \varepsilon \vartheta \in H_{g}$.

## Proof:

Utilizing Lemma 2.5 in [1] it suffices to prove that $H_{s} \cup H_{0} \subset H_{\pi}$, where $H_{\pi}$ is the subset of $v \in T^{2}$ such that $v_{1}-v_{2} \leqslant \pi$. By the previous theorem it is then enough to show that $H_{b} \subset H_{\pi}$. By definition of $A_{1}$ and $A_{2}, A_{1} \subset H_{\pi}, A_{2} \subset H_{\pi}$. Suppose that $v \in A_{3}$. Then $v_{1}\left|\leq \frac{\pi}{2}, w_{2}\right| \leq \frac{\pi}{2}$, implying $\left|v_{1} \rightarrow \vartheta_{2}\right| \leq \pi$, so that $A_{3} \subset H_{\pi}$ and the proof is complete.

Corollary 2.1 implies that $H_{B}$ is connected, that is, the stable region has only one component in contrast with the general case (see [1, Example 1]). Also, $H_{B} \cup H_{0}=\bar{H}_{3}$, where the bar denotes closure. Hence $H_{0}=\partial H_{s}$ (boundary of $\mathrm{H}_{\mathrm{a}}$ ). The next result characterizes this boundary.

Theorem 2.2 $\mathrm{H}_{0}=\partial \mathrm{H}_{\mathrm{a}}$ is a connected one dimensional differentiable manifold in $\mathrm{T}^{2}$.

Proof:
Connectedness follows easily from the previous discussion. To establish that $\partial \mathrm{H}_{\mathrm{a}}$ is a manifold, observe that

$$
\begin{equation*}
\partial \mathrm{H}_{\mathrm{a}}=\left\{\vartheta \in \mathrm{H}_{\mathrm{b}} \mid \operatorname{det} \mathrm{F}(v)=0\right\} \tag{2.5}
\end{equation*}
$$

Hence, it suffices to show that $\nabla_{*} \operatorname{det} \mathrm{~F}(v) \neq 0$ on $\partial \mathrm{H}_{\mathrm{s}}$. For $v \in \mathrm{~T}^{2}$ define the following real functions:

$$
\begin{align*}
& \psi_{1}(v)=[1,0] F(v)\left[\begin{array}{l}
1 \\
0
\end{array}\right]=B_{1} \cos v_{1}+B_{12} \cos \left(v_{1}-v_{2}\right)  \tag{2.6a}\\
& \psi_{2}(v)=[0,1] F(v)\left[\begin{array}{l}
0 \\
1
\end{array}\right]=B_{2} \cos v_{2}+B_{12} \cos \left(v_{1}-v_{2}\right)  \tag{2.6b}\\
& \psi_{3}(v)=[1,1] F(v)\left[\begin{array}{l}
1 \\
1
\end{array}\right]=B_{1} \cos v_{1}+B_{2} \cos v_{2} \tag{2.6c}
\end{align*}
$$

It is clear that these functions are nonnegative on $H_{0}$. Also note that $F(\vartheta)$ never vanishes on $\mathrm{T}^{2}$ since this would imply $\cos v_{1}=0, \cos v_{2}=0$, $\cos \left(v_{1}-v_{2}\right)=0$, an obvious contradiction. Hence, since $[1,0],[0,1]$ and $[1,1]$ are pairwise independent, it follows that at most one of the functions $\psi_{i}$ may vanish at any $v \in T^{2}$. Assume now that $v \in \partial H_{s}$ but $\nabla_{*} \operatorname{det} F(v)=0$. Without loss of generality we can suppose that $\psi_{2}(v)>0, \psi_{3}(v)>0$.

We have

$$
\nabla_{\theta} \operatorname{det} F(v)=-\left[\begin{array}{l}
\mathrm{B}_{1} \sin v_{1} \psi_{2}(v)+\mathrm{B}_{12} \sin \left(v_{1}-v_{2} \psi_{9}(v)\right.  \tag{2.7}\\
\mathrm{B}_{2} \sin v_{2} \psi_{1}(v)-\mathrm{B}_{12} \sin \left(v_{1}-v_{2}\right) \psi_{3}(v)
\end{array}\right]
$$

Therefore, $\nabla_{*}$ det $F(\vartheta)=0$ implies

$$
\begin{align*}
& \mathrm{B}_{1} \sin v_{1} \psi_{2}(v)+\mathrm{B}_{12} \sin \left(v_{1}-v_{2}\right) \psi_{3}(v)=0  \tag{2.8a}\\
& \mathrm{~B}_{2} \sin v_{2} \psi_{1}(v)-\mathrm{B}_{12} \sin \left(v_{1}-v_{2}\right) \psi_{3}(v)=0 \tag{2.8b}
\end{align*}
$$

Without loss of generality, assume that $\sin v_{1} \geq 0$. From (2.8a) we then have $\sin \left(v_{1}-v_{2}\right) \leq 0$. Since $v \in \partial H_{3} \leq H_{\pi}$ it follows that $0 \leq v_{1} \leq \pi,-\pi \leq v_{1}-v_{2} \leq 0$; thus, $0 \leq v_{2} \leq \pi$ or $\sin v_{2} \geq 0$. This contradicts (2.8) unless $\sin v_{1}=0$ and $\sin \left(v_{1}-v_{2}\right)=0$ in which case $v_{i}$ is a multiple of $\pi$. By Lemma 3 of [1] it is unstable and so $\vartheta$ cannot belong to $\partial \mathrm{H}_{\mathbf{s}}$. The proof is complete.

### 2.3 On the Convexity of $\mathrm{H}_{3}$.

It was conjectured in [2] that the set of stable solutions $H_{s}$ is convex. If this were true, one could conclude almost immediately that the stable solution of the load flow function is unique. Here we give a counterexample to this conjecture.

Theorem $2.3 \mathrm{H}_{8}$ is not generally convex.
Proof: Suppose $\mathrm{H}_{\mathrm{B}}$ is convex. Let $v \in \partial \mathrm{H}_{\mathrm{g}}$ be arbitrary. Let $\mathrm{v} \in \mathrm{T}_{\theta} \partial \mathrm{H}_{\mathrm{s}}$, the tangent space of $\partial \mathrm{H}_{s}$ at $v$. Then $v+t v, t \in R$, is a hyperplane (which is a straight line in this context) through $v$ supporting $H_{s}$. In addition, since $F(\vartheta)$ has one strictly positive eigenvalue for $v$ in $\partial \mathrm{H}_{\mathrm{s}}$ we conclude that for some $\varepsilon>0$, $\operatorname{det} F(v+t v) \leq 0$ for all $|t|<\varepsilon$. So it must be the case that

$$
\begin{equation*}
\mathbf{v}^{\mathrm{T}} \frac{\partial^{2}}{\partial \vartheta^{2}}(\operatorname{det} F(\vartheta)) \mathrm{v} \leq 0 . \tag{2.9}
\end{equation*}
$$

We can find explicit expressions for $v$ and $\frac{\partial^{2}}{\partial \vartheta^{2}}(\operatorname{det} F(\vartheta))$.

$$
v=\binom{v_{1}}{v_{2}}, \frac{\partial^{2}}{\partial \vartheta^{2}}(\operatorname{det} F(\vartheta))=\left[\begin{array}{ll}
a_{11} & a_{12} \\
a_{21} & a_{22}
\end{array}\right]
$$

where

$$
\begin{align*}
\mathrm{v}_{1} & =\mathrm{B}_{12} \sin \left(v_{1}-v_{2}\right) \psi_{3}(v)-\mathrm{B}_{2} \sin v_{2} \psi_{1}(v)  \tag{2.10a}\\
\mathrm{v}_{2} & =\mathrm{B}_{12} \sin \left(v_{1}-v_{2}\right) \psi_{3}(v)-\mathrm{B}_{1} \sin v_{1} \psi_{2}(v)  \tag{2.10b}\\
\mathrm{a}_{11} & =-\mathrm{B}_{1} \cos v_{1} \mathrm{~B}_{12} \cos \left(v_{1}-v_{2}\right)+2 \mathrm{~B}_{1} \sin v_{1} \mathrm{~B}_{12} \sin \left(v_{1}-v_{2}\right)  \tag{2.10c}\\
\mathrm{a}_{22} & =-\mathrm{B}_{2} \cos v_{2} \mathrm{~B}_{12} \cos \left(v_{1}-v_{2}\right)+2 \mathrm{~B}_{2} \sin v_{2} \mathrm{~B}_{12} \sin \left(v_{1}-v_{2}\right)  \tag{2.10d}\\
\mathrm{a}_{12} & =\mathrm{a}_{21}=-\mathrm{B}_{1} \cos v_{1} \mathrm{~B}_{2} \cos v_{2}+\mathrm{B}_{1} \sin v_{1} \mathrm{~B}_{2} \sin v_{2} \\
& +\mathrm{B}_{2} \sin v_{2} \mathrm{~B}_{12} \sin \left(v_{1}-v_{2}\right)-\mathrm{B}_{1} \sin v_{1} \mathrm{~B}_{12} \sin \left(v_{1}-v_{2}\right) \tag{2.10e}
\end{align*}
$$

Now let $v=\left(v_{1}, v_{2}\right)$ be given by $v_{1}=75^{\circ}, v_{2}=-60^{\circ}$. Consider the following sequence of parameter values $B^{n}=\left\{B_{1}^{n}, B_{2}^{n}, B_{12}^{n}\right\}$ given by

$$
B_{1}^{n}=2^{n}\left[\frac{\sqrt{2}}{2}+\frac{1}{2^{n}}\right] \frac{\sqrt{2}}{4}, B_{2}^{n}=\frac{\sqrt{2}}{2}+\frac{1}{2^{n}}, B_{12}^{n}=\frac{1}{2} .
$$

One can check that $\operatorname{det} F^{n}(v)=0$ for all $n$, where $F^{n}(v)$ denotes the Jacobian of the load flow function with parameters $B^{n}$. Furthermore, since $\cos v_{1}>0, \cos v_{2}>0$ we deduce that $\mathrm{Fr}(v) \geq 0$ for all $n$ and hence $v \in \partial \mathrm{H}_{\mathrm{g}}$. Denoting by $v_{i}^{n}, a_{i j}^{n}$ the values of $v_{i}, a_{i j}$ (see 2.10) at $B^{n}, v$ we investigate the asymptotic behavior of $\nabla_{i}^{n}, a_{i j}^{n}$ noting that terms containing $B_{1}^{n} \sim 2^{n-2}$ will dominate over the rest.

This leads to

$$
\begin{align*}
& \mathrm{v}_{1}^{\mathrm{n}} \sim-2^{\mathrm{n}-2}\left[\frac{\sqrt{2}}{4}(1+\sqrt{3})\right] \cos v_{1}  \tag{2.11a}\\
& \mathrm{v}_{2}^{\mathrm{n}} \sim-2^{\mathrm{n}-2} \frac{\sqrt{2}}{4} \cos v_{1}  \tag{2.11b}\\
& \mathrm{a}_{11}^{\mathrm{n}} \sim 2^{\mathrm{n}-2} \frac{\sqrt{2}}{4} \cos v_{1}+2^{\mathrm{n}-1} \frac{\sqrt{2}}{4} \sin v_{1}  \tag{2.11c}\\
& \mathrm{a}_{12}^{\mathrm{n}}=\mathrm{a}_{21}^{\mathrm{n}} \sim 2^{\mathrm{n}-2} \frac{\sqrt{2}}{4} \cos v_{1}-2^{\mathrm{n}-2} \frac{\sqrt{2}}{4}(a+\sqrt{3}) \sin v_{1} \tag{2.11d}
\end{align*}
$$

while $a_{2}^{n}$ is bounded above. Thus:

$$
\left(v^{n}\right)^{T} \frac{\partial^{2}}{\partial v^{2}}\left(\operatorname{det} F^{n}(v)\right) v^{n}=\left(v_{1}^{n}\right)^{2} a_{11}^{n}+2 v_{1}^{n} v_{2}^{n} a_{12}^{n}+\left(v_{2}^{n}\right)^{2} a_{2 R}^{n} \sim 2^{n-6} \sqrt{2},
$$

so that for $n$ large enough,

$$
\left(v^{n}\right)^{T} \frac{\partial^{2}}{\partial v^{2}}\left(\operatorname{det} F^{n}(v)\right) v^{n}>0
$$

contrary to the original hypothesis, and the proof is complete.

### 2.4 On the number of stable solutions

We now seek to establish the uniqueness of the stable solution of the load flow equation. The proof is rather involved, and for reasons of clarity, we first establish two lemmas.

Lemma 2.1 Let $v, \forall \in H_{s}$ be such that $v \neq \mathfrak{v}, f(v)=f(v)$. Then there exist $\xi \in \bar{H}_{s}$ and $\boldsymbol{\xi} \in \partial \mathrm{H}_{\mathrm{a}}$ such that $\mathrm{f}(\xi)=\mathrm{f}(\boldsymbol{\xi})$.

Proof: Since $v, \vartheta \in \mathrm{H}_{\mathrm{g}}, \operatorname{det} \mathrm{F}(\vartheta) \neq 0, \operatorname{det} \mathrm{~F}(\mathcal{v}) \neq 0$, and in view of the inverse function theorem, there exist disjoint neighborhoods $V, \nabla$ of $v$ and $v$ respectively, and a diffeomorphism $\eta: V \rightarrow \nabla$ such that if $x \in V$, then $f(x)=f(\eta(x)), x \neq \eta(x)$. Therefore, given the differentiable curve $\gamma(t)=(1-t) \vartheta$, there exists an open interval $[0, \alpha)$ and a differentiable curve $\tilde{\gamma}(t)$ defined on the interval such that $\tilde{\gamma}(0)=\mathcal{F}$ and $\mathrm{f}(\boldsymbol{\gamma}(\mathrm{t}))=\mathrm{f}(\gamma(\mathrm{t})), \boldsymbol{\gamma}(\mathrm{t}) \neq \gamma(\mathrm{t}), \tilde{\gamma}(\mathrm{t}) \in \mathrm{H}_{\mathrm{g}}$. Now let $\alpha \in[0,1]$ be the largest such real number. Let $\boldsymbol{\xi}=\lim _{t \rightarrow \alpha} \tilde{\gamma}(t)$ and let $\xi=(1-\alpha) v$. Observe first that $\boldsymbol{\xi} \neq \boldsymbol{\xi}$ since $(1-\alpha) \vartheta \in H_{s}$ by Corollary 2.1 and hence $F((1-\alpha) \vartheta)$ is nonsingular. Next
note that $\alpha<1$ because if $\alpha=1$ then $f(\underset{\xi}{ })=f(0)$ for some $\boldsymbol{\xi} \in \bar{H}_{8}$. But then by Corollary 2.1, $\varepsilon \boldsymbol{\xi} \in H_{a}$ for $\varepsilon \in(0,1)$. Hence,

$$
\left.0=\langle\dot{\xi}, \mathrm{f}(\xi)-\mathrm{f}(0)\rangle=\int_{0}^{1}\langle\dot{\xi}, F(\varepsilon \xi) \xi\rangle \mathrm{d} \varepsilon\right\rangle 0,
$$

which is a contradiction. It remains to show that $\boldsymbol{\xi} \in \partial \mathrm{H}_{\mathrm{g}}$. This is clear, for if家 $\in H_{s}$ then the definition of $\tilde{\gamma}(t)$ can be extended beyond $\alpha$, contradicting its maximality. The proof is complete.

Lemma 2.2 Let $\vartheta^{0} \in \partial \mathrm{H}_{\mathrm{g}} \cap \mathrm{A}$ where

$$
\mathrm{A}=\left\{v \in \mathrm{~T}^{2} \left\lvert\, \frac{\pi}{2} \leq v_{1} \leq \pi\right., 0 \leq v_{2} \leq \frac{\pi}{2}, 0 \leq v_{1}-v_{2} \leq \frac{\pi}{2}\right\} .
$$

Then,
$\mathrm{B}_{1} \sin v_{1}^{0}+\mathrm{B}_{2} \sin v_{2}^{0} \geq \mathrm{B}_{1}$,
$B_{1} \sin v_{1}^{0}+B_{12} \sin \left(v_{1}^{0}-v_{2}^{0}\right) \geq B_{1}$.
Proof: $\partial \mathrm{H}_{B}$ is a closed connected one dimensional manifold. Furthermore, $\left(\frac{\pi}{2}, \frac{\pi}{2}\right),\left(\frac{\pi}{2}, 0\right)$ are the only points in $\partial \mathrm{A} \cap \partial \mathrm{H}_{s}$. This assertion can be verified very easily. One can conclude, then, that there exists a differentiable curve $\gamma(\mathrm{t}): \mathrm{R} \rightarrow \mathrm{A}$ such that $\gamma(0)=\left(\frac{\pi}{2}, \frac{\pi}{2}\right), \gamma(1)=\left(\frac{\pi}{2}, 0\right), \quad \gamma\left(\mathrm{t}_{0}\right)=\vartheta^{0}, \quad 0 \leq \mathrm{t}_{0} \leq 1$, and $\gamma(\mathrm{t}) \in \partial \mathrm{H}_{\mathrm{a}} \cap \mathrm{A}$ for all $\mathrm{t} \in[0,1]$. Consider the function $\varphi: \mathrm{T}^{2} \rightarrow \mathrm{R}$ given by

$$
\begin{equation*}
\varphi(v)=\mathrm{B}_{1} \sin v_{1}+\mathrm{B}_{2} \sin v_{2} \tag{2.12}
\end{equation*}
$$

We have

$$
\begin{equation*}
\varphi(\gamma(0))=B_{1}+B_{2}, \varphi(\gamma(1))=B_{1} \tag{2.13}
\end{equation*}
$$

Hence it suffices to show that $\varphi(\gamma(\mathrm{t})$ ) is monotone on ( 0,1 ). Consider

$$
\begin{equation*}
\frac{d}{d t} \varphi(\gamma(t))=(1,1) F(\gamma(t)) \dot{\gamma}(t) \tag{2.14}
\end{equation*}
$$

But $\dot{\gamma}(\mathrm{t}) \in \mathrm{T}_{\gamma(\mathrm{t})}\left(\partial \mathrm{H}_{\mathrm{g}}\right)$, and $\mathrm{T}_{\mathrm{q}}\left(\partial \mathrm{H}_{\mathrm{g}}\right)$ is spanned by the vector $\mathrm{v}(\vartheta)$ given in (2.10). By direct computation,

$$
\begin{equation*}
(1,1) F(v) v(v)=-B_{2} \sin v_{2} B_{1} \cos v_{1} \psi_{1}(v) \tag{2.15}
\end{equation*}
$$

$$
+\mathrm{B}_{1} \sin v_{1} \mathrm{~B}_{2} \cos v_{2} \psi_{2}(v)+\mathrm{B}_{12} \sin \left(v_{1}-v_{2}\right)\left(\psi_{9}(v)\right)^{2}
$$

Finally, by inspection, (1,1) $\mathrm{F}(\vartheta) \mathrm{v}(\vartheta)>0$ on $\partial \mathrm{H}_{\mathrm{B}} \cap \operatorname{int}(\mathrm{A})$ and the proof is complete. (The proof of the second inequality follows by symmetry upon interchanging $v_{1}-v_{2}$ and $v_{2}$.)

We are now ready to prove the main result.
Theorem 2.4 The function f is one-to-one on $\overline{\mathrm{H}}_{\mathrm{g}}$.
Proof: By contradiction. Suppose there exist $v, \vartheta$ in $\bar{H}_{s}$ with $v \neq v$ and $f(\vartheta)=f(\vartheta)$. By Lemma 2.1 we can assume $\vartheta \in \partial H_{s}$, and in view of the symmetry of the load flow function we may assume $\tilde{v}_{1} \geq \frac{\pi}{2}$. Note that since $\forall \in \bar{H}_{8}$ we must necessarily have $0 \leq v_{2} \leq \frac{\pi}{2}, 0 \leq \mathcal{v}_{1}-\vartheta_{2} \leq \frac{\pi}{2}$ and Lemma 2.2 implies

$$
\begin{align*}
& f_{1}(\vartheta)=B_{1} \sin \left(\vartheta_{1}\right)+B_{12} \sin \left(\vartheta_{1}-\vartheta_{2}\right) \geq B_{1}  \tag{2.16a}\\
& f_{1}(\vartheta)+f_{2}(\mho)=B_{1} \sin \left(\vartheta_{1}\right)+B_{2} \sin \left(\vartheta_{2}\right) \geq B_{1} . \tag{2.16b}
\end{align*}
$$

Since $f(\vartheta)=f(\vartheta)$ we get

$$
\begin{align*}
& f_{1}(v) \geq B_{1}  \tag{2.17a}\\
& f_{1}(v)+f_{2}(v) \geq B_{1} \tag{2.17b}
\end{align*}
$$

Observe that (2.17) implies $\sin v_{2} \geq 0, \sin \left(v_{1}-v_{2}\right) \geq 0$, hence, since $v \in \bar{H}_{s}$ we must have

$$
0 \leq v_{1} \leq \pi, 0 \leq v_{2} \leq \frac{\pi}{2}, 0 \leq v_{1}-v_{2} \leq \frac{\pi}{2} .
$$

Now distinguish two cases:
i) $\sin \vartheta_{1} \geq \sin \vartheta_{1}$ :

This clearly implies

$$
\begin{align*}
& 0 \leq \pi-v_{1} \leq v_{1} \leq v_{1} \leq \pi  \tag{2.18a}\\
& 0 \leq v_{2} \leq v_{2} \leq \frac{\pi}{2}  \tag{2.18b}\\
& 0 \leq v_{1}-v_{2} \leq v_{1}-v_{2} \leq \frac{\pi}{2} \tag{2.18c}
\end{align*}
$$

Let $u \in R^{2}$ be given by

$$
u=\left[\begin{array}{l}
\tilde{\tau}_{1}-v_{1}  \tag{2.20}\\
\tilde{v}_{2}-v_{2}
\end{array}\right]
$$

By direct computation,

$$
\begin{aligned}
& \langle u, F(v+u t) u\rangle=B_{1} \cos \left(v_{1}+\left(\vartheta_{1}-v_{1}\right) t\right)\left(v_{1}-v_{1}\right)^{2} \\
& +\mathrm{B}_{2} \cos \left(\vartheta_{2}+\left(\vartheta_{2}-\vartheta_{2}\right) \mathrm{t}\right)\left(\vartheta_{2}-\vartheta_{2}\right)^{2} \\
& +\mathrm{B}_{12} \cos \left(\vartheta_{1}-\vartheta_{2}+\left(\vartheta_{1}-v_{2}-v_{1}+\vartheta_{2}\right) \mathrm{t}\right)\left(v_{1}-\vartheta_{2}-v_{1}+\vartheta_{2}\right)^{2} \\
& >\mathrm{B}_{1} \cos \left(\vartheta_{1}\right)\left(\mho_{1}-\vartheta_{1}\right)^{2}+\mathrm{B}_{2} \cos \tilde{\vartheta}_{2}\left(\psi_{2}-\vartheta_{2}\right)^{2} \\
& +B_{12} \cos \left(\vartheta_{1}-\vartheta_{2}\right)\left(\tilde{v}_{1}-\vartheta_{2}-v_{1}+v_{2}\right)^{2}=\left\langle u, F\left(v_{)}\right) u\right\rangle \geq 0
\end{aligned}
$$

for all $t$ in ( 0,1 ), and hence,

$$
\begin{equation*}
u^{T}(f(\vartheta)-f(v))=\int_{0}^{t}\langle u, F(v+u t) u\rangle d t>0 \tag{2.21}
\end{equation*}
$$

contradicting our initial hypothesis that $f(\widetilde{v})=f(\vartheta)$.
ii) $\sin \vartheta_{1}>\sin \vartheta_{1}$ :
$\dot{A}$ contradiction here is established in a similar manner and completes the proof.

It is clear that $f\left(H_{s}\right)$ is an open set in $R^{2}$. By the continuity of $f$ we also get $\partial f\left(H_{s}\right) \subset f\left(\partial H_{s}\right)$. In fact, we can now also establish that $\partial f\left(H_{8}\right)=f\left(\partial H_{8}\right)$. If this is not the case, then there exists $\vartheta^{0}$ in $\partial \mathrm{H}_{\mathrm{B}}$ such that $\mathrm{f}\left(\vartheta^{0}\right) \in \mathrm{f}\left(\mathrm{H}_{8}\right)$, in violation of Theorem 2.4.

### 2.5 On the range of $f$

We seek to characterize the range of the load flow function. We first establish that $f\left(H_{s}\right)$ is a convex set in $R^{2}$. The following lemma is useful.

Lemma 2.3 Let C be an open bounded set in $\mathrm{R}^{\mathrm{n}}$ such that $\partial \mathrm{C}$ is an ( $\mathrm{n}-1$ ) dimensional manifold. Suppose that for any three distinct points $x, y, z$, in $\partial C$, $T_{\mathbf{x}}(\partial C)=T_{y}(\partial C)$, implies $T_{8}(\partial C) \neq T_{\mathbf{x}}(\partial C)$. Then $C$ is convex.

Proof: By contradiction. Suppose C satisfies the hypothesis of the lemma and C is not convex. Then there exist $x_{1}$ and $x_{2}$ in $C$ such that $\not / 2\left(x_{1}+x_{2}\right)$ is not in $C$. Define $g: R \rightarrow R^{n}$ by

$$
\begin{equation*}
g(t)=t x_{1}+(1-t) x_{2} \tag{2.22}
\end{equation*}
$$

Let $t_{0}$ in $(0,1)$ be such that $g\left(t_{0}\right) \in \partial C$. Let $x_{0}=g\left(t_{0}\right)$ and consider $T_{x_{0}}(\partial C)$ and a vector $w_{0}$ normal to $T_{x_{0}}(\partial C)$. Define the linear functional $\lambda$ on $C$ by

$$
\begin{equation*}
\lambda(x)=\left\langle x, w_{0}\right\rangle \tag{2.23}
\end{equation*}
$$

Since $C$ is bounded, $C$ is compact, there exist points $x_{m p} x_{M}$ where $\lambda$ attains a global minimum and maximum respectively.

Obviously, $x_{m} \neq x_{M}$. Furthermore, since

$$
\lambda\left(x_{0}\right)=t_{0} \lambda\left(x_{1}\right)+\left(1-t_{0}\right) \lambda\left(x_{2}\right)
$$

and $x_{1}, x_{2}$ are in $C$, it is evident that $x_{0} \neq x_{m}$, and $x_{0} \neq x_{M}$, while

$$
\mathrm{T}_{\mathbf{x}_{\mathbf{m}}}(\partial \mathrm{C})=\mathrm{T}_{\mathbf{x}_{\mathbf{w}}}(\partial \mathrm{C})=\mathrm{T}_{\mathbf{x}_{0}}(\partial \mathrm{C})
$$

contrary to the hypothesis, and the proof is complete.
Lemma $2.4 \partial f\left(\mathrm{H}_{3}\right)=f\left(\partial \mathrm{H}_{3}\right)$ is a one-dimensional differentiable manifold in $\mathrm{R}^{2}$.
Proof: Note that $\mathrm{f}: \partial \mathrm{H}_{\mathrm{s}} \rightarrow \partial \mathrm{f}\left(\mathrm{H}_{\mathrm{s}}\right)$ is a continuous bijective map. Thus it suffices to prove that $F(v)$ is bijective on $\mathrm{T}_{v}\left(\partial \mathrm{H}_{8}\right)$ for all $v$ in $\partial \mathrm{H}_{8}$. Without loss of generality, assume $v \in \mathrm{~A}$ (see Lemma 2.2), i.e., $\frac{\pi}{2} \leq v_{1} \leq \pi, 0 \leq v_{2} \leq \frac{\pi}{2}, 0 \leq v_{1}-v_{2} \leq \frac{\pi}{2}$.

We already know from the proof of Lemma 2.2 that $F(v) v(v) \neq 0$ for all $v$ in $\partial \mathrm{H}_{3} \cap \operatorname{Int}(\mathrm{~A})$ where $v(v)$ is the vector defined in (2.10) which spans $\mathrm{T}_{8}\left(\partial \mathrm{H}_{8}\right)$. So it remains to show that $F(v) v(v) \neq 0$ on $\partial H_{s} \cap \partial A$, namely, at $\left(\frac{\pi}{2}, \frac{\pi}{2}\right)$ and $\left(\frac{\pi}{2}, 0\right)$.

At $\left(\frac{\pi}{2}, \frac{\pi}{2}\right)$ we get

$$
F(v) v(v)=\left[\begin{array}{cc}
B_{12} & -B_{12} \\
-B_{12} & B_{12}
\end{array}\right]\left[\begin{array}{c}
-B_{2} B_{12} \\
B_{1} B_{12}
\end{array}\right]=\left(B_{1}+B_{2}\right) B_{12}^{2}\left[\begin{array}{c}
-1 \\
1
\end{array}\right] \neq 0 .
$$

On the other hand at ( $\frac{\pi}{2}, 0$ ) the assertion follows by inspection from (2.15) and the proof is complete.

Theorem $2.5 \mathrm{f}\left(\mathrm{H}_{\mathrm{g}}\right)$ is an open convex set in $\mathrm{R}^{\text {2 }}$.

Proof: Observe that if $\approx$ belongs to $\partial \mathrm{H}_{s}$, then so does $-\vec{v}$ and $\mathrm{T}_{\mathrm{f}(\tilde{v})}\left(\partial \mathrm{f}\left(\mathrm{H}_{8}\right)\right)=$ $T_{f(-g)}\left(\partial f\left(H_{s}\right)\right)$ so that in view of the lemmas above, it suffices to show that there exists no $v \in \partial \mathrm{H}_{\mathrm{s}}$ such that $v \neq v, v \neq \cdots$ and $\mathrm{F}(v)=\alpha F(\vartheta)$ for some $\alpha \neq 0$.

We proceed by contradiction. Let $v, \vartheta$ be as above. Without loss of generality we can assume $\mathfrak{v} \in \mathrm{A}$, i.e., $\frac{\pi}{2} \leq \mathcal{v}_{1} \leq \pi, 0 \leq \tilde{v}_{2} \leq \frac{\pi}{2}, 0 \leq \mathcal{v}_{1}-\mathcal{v}_{2} \leq \frac{\pi}{2}$. Evidently, we must have $\alpha>0$, so again without loss of generality we can assume $\alpha \geq 1$, as well as $v_{1} \geq 0$. Then from $F(v)=\alpha F(v)$ we obtain

$$
\begin{align*}
& \cos v_{1}=\alpha \cos v_{1}  \tag{2.24a}\\
& \cos v_{2}=\alpha \cos v_{2}  \tag{2.24b}\\
& \cos \left(v_{1}-v_{2}\right)=\alpha \cos \left(v_{1}-v_{2}\right) \tag{2.24c}
\end{align*}
$$

Hence, in view of our assumptions,

$$
\begin{align*}
& \cos v_{1} \leq \cos v_{1} \leq 0  \tag{2.25a}\\
& \cos v_{2} \geq \cos v_{2} \geq 0  \tag{2.25b}\\
& \cos \left(v_{1}-v_{2}\right) \geq \cos \left(v_{1}-w_{2}\right) \geq 0 \tag{2.25c}
\end{align*}
$$

Since $v_{1} \geq 0$, we obtain from (2.25a)

$$
\begin{equation*}
\frac{\pi}{2} \leq v_{1} \leq v_{1} \leq \pi \tag{2.26a}
\end{equation*}
$$

And from (2.25b) and (2.25c)

$$
\begin{align*}
& v_{2} \leq v_{2} \leq \frac{\pi}{2}  \tag{2.26b}\\
& v_{1}-v_{2} \leq v_{1}-v_{2} \leq \frac{\pi}{2} \tag{2.26c}
\end{align*}
$$

But adding (2.26b) and (2.26c) gives $v_{1} \leq \mathfrak{v}_{1}$, in violation of (2.26a). The proof is complete.

So far we have established that $f\left(H_{s}\right)$ is a convex set. Notice that $f\left(H_{s}\right)$ is thus the convex hull of $f\left(\partial \mathrm{H}_{\mathrm{B}}\right)$. The question that arises now is whether $f\left(\bar{H}_{g}\right)=f\left(T^{2}\right)$; i.e., whether we can always guarantee the existence of a stable solution. Observe that the point $\left(\frac{\pi}{2}, \frac{\pi}{2}\right)$ is in $\partial \mathrm{H}_{3}$, and in addition, the function
$f_{1}(v)+f_{2}(v)$ attains a global maximum over $f\left(T^{2}\right)$ at $\left(\frac{\pi}{2}, \frac{\pi}{2}\right)$. Hence, $f\left(\frac{\pi}{2}, \frac{\pi}{2}\right)$ belongs to $\partial f\left(\mathrm{~T}^{2}\right)$ and $f\left(\partial \mathrm{H}_{8}\right) \cap \partial f\left(\mathrm{~T}^{2}\right)$ is nonempty.

After this initial observation we proceed to establish the theorem concluding this section.

Theorem $2.6 \mathrm{f}\left(\mathrm{H}_{\mathrm{s}}\right)=\mathrm{f}\left(\mathrm{T}^{2}\right)$
Proof: By contradiction. Suppose $f\left(\mathrm{H}_{\mathrm{g}}\right) \neq \mathrm{f}\left(\mathrm{T}^{2}\right)$. From the preceding paragraph,

$$
\begin{equation*}
\partial f\left(\bar{H}_{s}\right) \cap \partial f\left(\mathrm{~T}^{2}\right) \neq \phi \tag{2.28}
\end{equation*}
$$

On the other hand, since $f\left(H_{8}\right) \neq f\left(T^{2}\right)$, and $f\left(H_{g}\right)$ is convex, therefore

$$
\begin{equation*}
\partial \mathrm{f}\left(\mathrm{~T}^{2}\right)-\partial \mathrm{f}\left(\overline{\mathrm{H}}_{\mathrm{s}}\right) \neq \phi \tag{2.29}
\end{equation*}
$$

Let $H_{0_{-}}$denote the set of points in $T^{2}$ such that $F(v) \leq 0$ but $F(v)$ is not strictly negative definite. It is evident that $\partial f\left(\mathrm{~T}^{2}\right)-\partial f\left(\mathrm{H}_{\mathrm{g}}\right) \subset \mathrm{f}\left(\mathrm{H}_{0_{-}}\right)$. Also, because $\mathrm{f}\left(\mathrm{T}^{2}\right)$ is connected and $f\left(\bar{H}_{s}\right)$ is convex, it must be the case that

$$
\partial \mathrm{f}\left(\overline{\mathrm{H}}_{\mathrm{s}}\right) \cap \mathrm{f}\left(\mathrm{H}_{0_{-}}\right) \neq \phi
$$

Equivalently, since $\partial f\left(\bar{H}_{8}\right)=f\left(\partial H_{8}\right)=f\left(H_{0}\right)$, we must have
$f\left(\mathrm{H}_{0}\right) \cap \mathrm{f}\left(\mathrm{H}_{0_{-}}\right) \neq \phi$
The proof will, therefore, be complete if we show that (2.30) is false. Let $v \in H_{0-}, \vartheta \in H_{0}$ be arbitrary. It should be evident by now that there is no loss of generality if we assume that $\mathfrak{v} \in \operatorname{int}(A)$, where the set $A$ was defined in Lemma 2.2. In other words, we have

$$
\begin{equation*}
\frac{\pi}{2}<\vartheta_{1}<\pi, 0<\mathfrak{v}_{2}<\frac{\pi}{2}, 0<\mathfrak{v}_{1}-\vartheta_{2}<\frac{\pi}{2} \tag{2.31}
\end{equation*}
$$

Also, since $v \in H_{0-}, \forall \in H_{0}$, we obtain

$$
\begin{align*}
& \mathrm{B}_{1} \cos v_{1}+\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right) \leq 0 \leq \mathrm{B}_{1} \cos \vartheta_{1}+\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right)  \tag{2.32a}\\
& \mathrm{B}_{2} \cos v_{2}+\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right) \leq 0 \leq \mathrm{B}_{2} \cos \vartheta_{2}+\mathrm{B}_{12} \cos \left(v_{1}-v_{2}\right)  \tag{2.32b}\\
& \mathrm{B}_{1} \cos v_{1}+\mathrm{B}_{2} \cos \vartheta_{2} \leq 0 \leq \mathrm{B}_{1} \cos v_{1}+\mathrm{B}_{2} \cos \vartheta_{2} . \tag{2.32c}
\end{align*}
$$

Multiplying (2.32a) and (2.32b) and noticing that $\operatorname{det} F(\vartheta)=\operatorname{det} F(\vartheta)=0$ we get $\cos ^{2}\left(v_{1}-v_{2}\right) \leq \cos ^{2}\left(v_{1}-v_{2}\right)$. Doing the same thing with (2.32a) and (2.32c) as well
as with (2.32b) and (2.32c) we finally obtain

$$
\begin{equation*}
\cos ^{2} \vartheta_{1} \leq \cos ^{2} \vartheta_{1}, \cos ^{2} \vartheta_{2} \leq \cos ^{2} \vartheta_{2}, \cos ^{2}\left(\vartheta_{1}-v_{2}\right) \leq \cos ^{2}\left(\vartheta_{1}-v_{2}\right) \tag{2.33}
\end{equation*}
$$

It must be the case, then, that

$$
\begin{equation*}
\sin ^{2} \vartheta_{1} \geq \sin ^{2} \vartheta_{1}, \sin ^{2} \vartheta_{2} \geq \sin ^{2} \vartheta_{2}, \sin ^{2}\left(\vartheta_{1}-v_{2}\right) \geq \sin ^{2}\left(\vartheta_{1}-v_{2}\right) \tag{2.34}
\end{equation*}
$$

Now we utilize Lemma 2.2 to obtain
$B_{1} \sin \vartheta_{1}+B_{12} \sin \left(v_{1}-v_{2}\right)=B_{1} \sin \vartheta_{1}+B_{2} \sin \left(\vartheta_{1}-\vartheta_{2}\right)>B_{1}$
$B_{1} \sin \vartheta_{1}+B_{2} \sin \vartheta_{2}=B_{1} \sin \vartheta_{1}+B_{2} \sin \vartheta_{2}>B_{1}$
Hence, $\sin v_{2}>0, \sin \left(v_{1}-v_{2}\right)>0$. which combined with (2.34) yields
$\sin v_{2} \geq \sin v_{2}, \sin \left(v_{1}-v_{2}\right) \geq \sin \left(v_{1}-v_{2}\right)$
Hence, there are only two alternatives: either

$$
\sin v_{1} \geq \sin v_{1}
$$

which immediately yields

$$
v=v
$$

contradicting the hypothesis, or

$$
\sin v_{1} \leq-\sin \mho_{1}
$$

in which case, from (2.31) and (2.36),

$$
\begin{aligned}
& -\pi<-v_{1} \leq v_{1} \leq v_{1}-\pi<0 \\
& 0<\vartheta_{2} \leq v_{2} \leq \pi-v_{2}<\pi
\end{aligned}
$$

so that

$$
-2 \pi<-\pi-\left(v_{1}-v_{2}\right) \leq v_{1}-v_{2} \leq\left(v_{1}-v_{2}-\pi\right)<0
$$

which yields $\sin \left(v_{1}-v_{2}\right)<\sin \left(v_{1}-v_{2}\right)$ unless

$$
v_{1}=-v_{1}, v_{2}=\pi-v_{2}
$$

which in turn contradicts (2.35). The proof is complete.

## 3. Remarks on the Dynamic Behavior

### 3.1 The Differential Equation

According to the classical model, the motion of the 3-bus power network is governed by the differential equations

$$
\begin{align*}
& \mathrm{m}_{1} \ddot{v}_{1}+\mathrm{d}_{1} \dot{v}_{1}=\mathrm{P}_{1}-\mathrm{f}_{1}\left(v_{1}, v_{2}\right)  \tag{3.1}\\
& \mathrm{m}_{2} \ddot{v}_{2}+\mathrm{d}_{2} \dot{v}_{2}=\mathrm{P}_{2}-\mathrm{f}_{2}\left(v_{1}, v_{2}\right)
\end{align*}
$$

where the third bus is taken to be a slack bus and the vector $\mathrm{P}=\left(\mathrm{P}_{1}, \mathrm{P}_{2}\right)$ corresponds to the mechanical power input. It is not difficult to show analytically that (3.1) has at most six equilbrium points in $\mathrm{T}^{2}$ corresponding to the solutions of

$$
\begin{equation*}
P=f(v) \tag{3.2}
\end{equation*}
$$

Thus, if $P$ is a regular value of $f(v)$, i.e., $\operatorname{det} F(v) \neq 0$, the points in the inverse image $f^{-1}(P)$ must fall under one of the following cases (see Figure 3):
i) for $P$ in $R_{2}, f^{-1}(P)$ contains two points in $T^{2}$, one of which is stable and the other is a saddle with one positive eigenvalue.
ii). for $P$ in $R_{4}, f^{-1}(P)$ contains four points in $T^{2}$, one of which is stable, two are saddles with one positive eigenvalue, and the fourth a saddle with two positive eigenvalues; and finally
iii) for $P$ in $R_{6}, f^{-1}(P)$ contains six points in $T^{2}$, one of which is stable, three are saddle points with one positive eigenvalue, and the rest are saddles with two positive eigenvalues.

In the statements above the eigenvalues refer to the linearization of (3.1) around the equilibrium.

In the study of the dynamics of (3.1) we should expect to encounter saddle connection bifurcations giving rise to much more complex phenomena than in the one-dimensional problem studied in [3]. The next situation is devoted to the study of a certain kind of dynamic bifurcation.

### 3.2 The Role of Damping in Complete Stability

The motivation of this section stems from the study of the one-dimensional problem as well as the characterization of complete stability to be found in [3].

A dynamical system is said to be completely stable if every trajectory converges to an equilibrium.

Consider the one-dimensional classical equation for a power network:

$$
\begin{equation*}
\ddot{v}+\alpha \dot{v}=\beta-\sin v \tag{3.3}
\end{equation*}
$$

The global properties of (3.3) are well known (see [3,4]). In particular it is known that if $\beta<1$; i.e. as long as a stable equilibrium exists, then the differential system is completely stable provided that the damping constant $\alpha$ is large enough. This fact, combined with the role played by damping in the characterization of complete stability in [3], suggests the conjecture that sufficiently large damping should produce complete stability in higher dimensional problems as well. The following example shows that this is false.

Example 3.1 Consider the three node network of Figure 4. The differential equations for such a system are

$$
\begin{align*}
& \ddot{v}_{1}+d \dot{v}_{1}=\left(\varepsilon-\varepsilon^{2}\right)-\varepsilon \sin v_{1}-\sin \left(v_{1}-v_{2}\right)  \tag{3.4}\\
& \ddot{v}_{2}+d \dot{v}_{2}=\left(\varepsilon-\varepsilon^{2}\right)-\varepsilon \sin v_{2}-\sin \left(v_{1}-v_{2}\right)
\end{align*}
$$

Let $K_{n}$ be the closed subset of $R^{4}$ defined by

$$
\begin{equation*}
K_{n}=\left\{\left(\dot{v}_{1}, \dot{v}_{2}, v_{1}, v_{2}\right) \mid \dot{v}_{1}=\dot{v}_{2}, v_{1}-v_{2}=2 \pi \mathrm{n}\right\} \tag{3.5}
\end{equation*}
$$

The projections of the $K_{n}$ on the $\left(\dot{v}_{1}=0, \dot{v}_{2}=0\right)$ plane are shown in Figure 5. Notice that each $\mathrm{K}_{\mathrm{n}}$ is an invariant set of the motion of (3.4), and that if $\varepsilon<\frac{4}{5}$, then all the equilibrium points of (3.4) lie in the sets $\mathrm{K}_{\mathrm{n}}$.

By changing coordinates to $\mathrm{z}=v_{1}-v_{2}$, we obtain from (3.4):

$$
\begin{equation*}
\ddot{z}+d \dot{z}=-\varepsilon\left(\sin v_{1}-\sin v_{2}\right)-2 \sin z \tag{3.6}
\end{equation*}
$$

It is straightforward to show from (3.6) that there exist $\varepsilon_{0}>0, d_{0}>0$ such that the invariant sets $K_{n}$ are stable for all $\varepsilon<\varepsilon_{0}, d>d_{0}$. (A closed set is said to be stable if every open neighborhood of the set contains a positively invariant neighborhood.)

Let $A_{n}$ be the region of attraction of $K_{n}$. The $A_{n}$ are open disjoint sets of $R^{4}$ and hence their union does not cover $\mathrm{R}^{4}$. Thus the complement of this union is a nonempty invariant set and it contains no equilibrium points. Hence a trajectory which starts in this complement can never converge and so the system is not completely stable.

### 3.3 Complets Stability Revisited

In the example of section 3.2, we showed that arbitrarily large damping need not lead to complete stablity. Upon closer examination of equation (3.4) one can see that for $\varepsilon=1$ (3.4) not only has six equilibrium points in $T^{2} \times R^{2}$, but it is also completely stable for all positive values of the damping $d$. So, if we vary $\varepsilon$ continuously from $\varepsilon=1$ to $\varepsilon=0$, we should expect not only static bifurcations, i.e. appearances or disappearances of equilibrium points, to take place, but dynamic bifurcations, i.e. saddle connection bifurcations, as well. In Figures 6a and 6 b we depict the abstract flow diagrams of (3.4) for $\varepsilon=1$ and $\varepsilon=\frac{4}{5}$. At $\varepsilon=\frac{4}{5}$ the number of critical points changes from 6 to 2 . Notice how the stable manifolds and unstable manifolds fuse together.

The following proposition throws some light on the mechanism of the loss of complete stability.

Proposition 3.1 Consider a three node network such that there are exactly two equilibrium points. Then the system is not completely stable.

Proof: By contradiction. Suppose that the system is completely stable. Let $v^{s}=\left(\vartheta_{1}^{3}, \vartheta_{2}^{3}\right)$ and $\vartheta^{u}=\left(\vartheta_{1}^{u}, \vartheta_{2}^{u}\right)$ be the stable and unstable points respectively. Now, consider the system behavior in $R^{4}$. Let $W^{u}\left(\vartheta^{u}\right)$ be the 1-dimensional unstable manifold of $\vartheta^{u}$. By the assumption of complete stability, $W^{u}\left(\vartheta^{u}\right)$ must be bounded. Hence $\overline{W^{u}\left(\vartheta^{u}\right)}$ is a compact one dimensional manifold with boundary. Let $\vartheta+2 \pi(m, n)$ denote the point $\left(v_{1}+2 \pi m, v_{2}+2 \pi n\right)$. The boundary of
$\overline{W^{u}\left(\vartheta^{u}\right)}$ consists of critical points. It is not hard to show, hence we omit it here, that there is no loss of generality if we assume that

$$
\begin{equation*}
\partial \overline{W^{u}}\left(\vartheta^{u}\right)=\left\{\vartheta^{s}+2 \pi(\mathrm{~m}, \mathrm{n}), \vartheta^{\mathrm{s}}+2 \pi(\tilde{\mathrm{~m}}, \tilde{\mathrm{n}})\right\} \tag{3.7}
\end{equation*}
$$

for some pairs of integers ( $m, n$ ) and ( $\tilde{m}, \tilde{n}$ ).
Now, let $\vartheta_{\mathrm{k}, 1}^{u}$ be the point

$$
\begin{equation*}
v_{\mathrm{k}, 1}^{u}=v^{u}+2 \pi(\operatorname{lm}-k \tilde{m}, \ln -k \tilde{n}) \tag{3.8}
\end{equation*}
$$

for some integers $l$ and $k$.
Let

$$
\begin{equation*}
N_{k}=\bigcup_{1} \overline{W^{u}\left(\vartheta_{k, 1}^{u}\right)} \tag{3.9}
\end{equation*}
$$

The sets $N_{k}$ are closed, invariant and pairwise disjoint. In addition, each $N_{k}$ is stable, which completes the proof of the proposition.

We conclude with the following theorem which is stated without proof.
Theorem 3.1 Consider the ( $n+1$ )-node power network governed by the equation

$$
\begin{equation*}
\mathrm{M} \ddot{\vartheta}+\mathrm{D} \dot{\vartheta}=\mathrm{P}-\mathrm{f}(\hat{v}) \tag{3.10}
\end{equation*}
$$

where M and D are n -dimensional diagonal matrices. Let $\vartheta_{\mathrm{i}}, \mathrm{i}=1, \ldots, \mathrm{k}$ in $\mathrm{T}^{\mathrm{n}}$ be the equilibrium points. Then (3.10) is completely stable if and only if:
i) $W^{n}\left(\vartheta_{1}\right)$ is bounded in $R^{n} \times R^{n}, i=1, \ldots, k$.
ii) $\bigcup_{v \in \Pi}^{\bigcup} \bigcup_{i=1}^{\mathbf{k}} \overline{W^{4}\left(\vartheta_{i}+v\right)}$ is a connected set in $R^{n} \times R^{n}$ where $\Pi$ is the set of vectors $v$ all of whose components are multiples of $2 \pi$.

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## Figure Captions

Fig 1 The three node network
Fig 2 The sets $\mathrm{H}_{\mathrm{s}}, \mathrm{H}_{\mathrm{p}}$
Fig 3 Number of solutions of（3．2）
Fig 4 Network of Example 3.1
Fig 5 The sets $K_{n}$
Fig 6a Flow diagram for $\varepsilon=1$
Fig 6b Flow diagram for $\varepsilon=\frac{4}{5}$


Fig l. The three node network


Fig. 2 The sets $H_{p}, H_{b}$


Fig. 3


$$
\begin{array}{ll}
P_{i}=\left(\epsilon-\epsilon^{2}\right) & i=1,2 \\
M_{i}=1 & i=1,2 \\
d_{i}=d & i=1,2 \\
0<\epsilon<1 &
\end{array}
$$

Fig. 4 The network of example 3.1


Fig. 5


Fig. 6a: $\epsilon=1$


Fig. 6b: $\epsilon=4 / 5$


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