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Memorandum No. UCB/ERL M95/10
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# An Accurate Time Domain Interconnect Model of Transmission Line Networks 

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January 30, 1995


#### Abstract

In this paper. we present a new time domain model of interconnects modeled as transmission line networks. Each element of the characteristics of a transmission line is modeled by a principal part and a remainder. The principal part consists of an impulse and an exponential function, whose Laplace transform matches the original function at infinity frequency with order 1 and at zero frequency with order 0 . The remainder in the time domain consists of a cubic polynomial for a single line and a pieceuise cubic polynomial for coupled lines. The model is stable, accurate, simple and efficient to use.


## 1 Introduction

With the rapid increase of the signal frequency and decrease of the feature sizes in high speed electronic circuits, interconnects play increasingly important roles. On the MCM and PCB level, interconnects are modeled as transmission line networks.

Many papers have been published in the analysis of transmission line networks. As transmission lines are characterized in frequency domain and are generally terminated in nonlinear loads. a time domain model of transmission lines is needed so that conrolution can be used to find the transient response. When inverse Fourier transform

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is directly used to find the model [2], the computation complexity will be proportional to the square of the simulation time, which is not efficient. Padé approximation [7] and moment matching method [8] have been used to approximate the characteristics of the line in frequency domain and then the time domain model is formed by inverse Laplace transform. While efficient recursive convolution can be accomplished by using such a model, there is no guarantee of the stability of the model, and it is difficult to predict the order of matching so that the model is accurate enough. Optimization techniques are suggested to find a good approximation in frequency domain and then transfer it into time domain [9, 4], but optimization has its known limits and it is diffcult to guarantee a global rather than a local optimum unless the expensive simulated annealing is used.

In this paper, we suggest a new approach to form a simple and accurate time domain model for transmission lines. The approach consists of two steps. The first step is an approximation in the frequency domain. This is done by matching the characteristics at frequency zero with order zero (matching the value only) and at frequency $\infty$ with order one (matching both the value and the derivative w.r.t. $1 / \mathrm{s}$ at $s=\infty$ ). This approximation is designated as the principal part of the model as it approximates the original characteristic well in a wide range of high frequencies. This principal part is a first order rational function with its inverse Laplace transform consisting of an impulse and an exponential function. The main discrepancy of the approximation occurs at low frequencies. except frequency zero. The difference between the original characteristics and its principal part is called the remainder. The second step is finding the values of the remainder at sampling frequencies of interest, using the inverse FFT to find the time domain function, and fitting it with a cubic polynomial or a piecewise cubic polynomial. Thus. the time domain model of each characteristic of a line consists of an impulse, an exponential function and a number of cubic polynomials. The model is stable, simple and accurate. As recursive convolution can be easily implemented with such a model, this model is efficient to use.

This paper is organized as follows. The model of a single transmission line is presented in Sec.2. The model of coupled transmission line is presented in Sec.3. The technique of piecewise cubic fitting and the formulas of recursive convolution with cu-
bic polynomial are presented in Sec. 4 and 5, respectively. Examples and conclusions are given in Sec.6.

## 2 Model of single uniform line

Let $r, l, c, g$ and $d$ be the resistance, inductance, capacitance and conductance per unit length, and the length of a line, respectively. The telegrapher's equations of the line is as follows:

$$
\begin{align*}
& \frac{d V(x, s)}{d x}=-Z(s) I(x, s)  \tag{1}\\
& \frac{d I(x, s)}{d x}=-Y(s) V(x, s) \tag{2}
\end{align*}
$$

where $Z(s)=r+s l . Y(s)=g+s c$, and $x=0$ and $x=d$ correspond to the near end and far end of the line, respectively.

A transmission line can be modeled by its characteristic 2 port [15] as shown in Fig.1, where $I_{1}(s)=V(0, s), I_{1}(s)=I(0, s), V_{2}(s)=V(d, s)$ and $I_{2}(s)=I(d, s)$. The 2 -port elements are the following:

$$
\begin{equation*}
Z_{\mathrm{c}}(s)=\sqrt{\frac{s l+r}{s c+g}} \tag{3}
\end{equation*}
$$

is the characteristic impedance, and

$$
\begin{align*}
& W_{1}(s)=\Gamma(s)\left(2 V_{2}(s)-W_{2}(s)\right)  \tag{4}\\
& W_{2}(s)=\Gamma(s)\left(2 W_{1}(s)-W_{1}(s)\right) \tag{5}
\end{align*}
$$

where $\Gamma(s)$ is the propagation function, which is equal to $\exp (-\theta(s) d)$ with

$$
\begin{equation*}
\theta(s)=\sqrt{(s l+r)(s c+g)} \tag{6}
\end{equation*}
$$

The characteristic 2-port can then be described by the following equations in frequency domain:

$$
\begin{equation*}
V_{1}(s)=Z_{c}(s) I_{1}(s)+W_{1}(s) \tag{7}
\end{equation*}
$$

and

$$
\begin{equation*}
V_{2}(s)=W_{2}(s)-Z_{c}(s) I_{2}(s) \tag{8}
\end{equation*}
$$

Let $v_{k}(t), i_{k}(t),(k=1,2), z_{c}(t)$ and $\gamma(t)$ be the inverse Laplace transform of $V_{k}(s)$, $I_{k}(s), Z_{c}(s)$ and $\Gamma(s)$, respectively: Then, the time domain characteristic 2-port model is described by the following equations:

$$
\begin{align*}
& v_{1}(t)=z_{c}(t) * i_{1}(t)+w_{1}(t)  \tag{9}\\
& v_{2}(t)=w_{2}(t)-z_{c}(t) * i_{2}(t) \tag{10}
\end{align*}
$$

with

$$
\begin{align*}
& w_{1}(t)=\gamma(t) *\left(2 v_{2}(t)-w_{2}(t)\right)  \tag{11}\\
& w_{2}(t)=\gamma(t) *\left(2 v_{1}(t)-w_{1}(t)\right) \tag{12}
\end{align*}
$$

where ${ }^{\text {* }}$ is the symbol for convolution.
The modeling of a single line consists of the modeling of its characteristic impedance $=c(t)$ and its propagation function $\gamma(t)$.

### 2.1 Model of characteristic impedance

Let $Z_{c}(s)$ be first approximated by $Z_{c p}(s)$, such that $Z_{c p}(0)=Z_{c}(0), Z_{c p}(\infty)=Z_{c}(\infty)$ and $Z_{c p}^{(-1)}=Z_{c}^{(-1)}$. where $F^{(-1)}=d F(s) /\left.d(1 / s)\right|_{0 \rightarrow \infty}{ }^{1} Z_{c}(0)=\sqrt{r / g}, Z_{c}(\infty)=\sqrt{l / c}$ and $Z_{c}^{(-1)}=\frac{1}{2} \sqrt{\frac{1}{c}}\left(\frac{r}{l}-\frac{g}{c}\right)$. Let

$$
\begin{equation*}
Z_{c p}(s)=\frac{a_{1} s+a_{0}}{s+b_{0}} \tag{13}
\end{equation*}
$$

Then. $a_{1}=Z_{c}(x)$.

$$
\begin{equation*}
b_{0}=\frac{Z_{c}^{(-1)}}{Z_{c}(0)-Z_{c}(\infty)} \tag{14}
\end{equation*}
$$

and $a_{0}=Z_{c}(0) b_{0}$. Let $z_{c p}(t)=L^{-1} Z_{c p}(s)$, then

$$
\begin{equation*}
z_{c p}(t)=Z_{c}(\infty) \delta(t)+k_{z} \exp \left(-b_{0} t\right) \tag{15}
\end{equation*}
$$

where $k_{z}=Z_{c}^{(-1)} . z_{c p}(t)$ is the principal part of $z_{c}(t)$.
Let $Z_{c r}(s)=Z_{c}(s)-Z_{c p}(s)$ be the remainder of $Z_{c}(s)$. Suppose that $f_{0}$ is the basic frequency, and $f_{\max }$ is the highest frequency of interest. When the source signal is a pulse with a rising and a falling time $t r$ and $t f$, respectively and assuming that

[^0]$t s$ is the simulation time, then we choose $f_{0}=1 / t s$ and $f_{\max }=2 / \min (t r, t f)$. Let $N^{r}=\min _{\mathrm{n}}\left\{2^{n} f_{0} \geq f_{\text {max }}\right\}$. Compute $Z_{\text {cr }}(s)$ at sampling frequencies $s_{k}=2 \pi k f_{0} j$ for k from 0 to $N-1$, and find $z_{c r}(t)=L^{-1} Z_{\text {cr }}(s)$ by inverse FFT. $z_{\text {cr }}(t)$ looks like a parabolic, and we use a cubic polynomial $z_{c r}^{\prime}(t)=a_{3} t^{3}+a_{2} t^{2}+a_{1} t+a_{0}$ to do least square fitting for $z_{\mathrm{r}}(t)$. Then, $z_{\mathrm{c}}(t)$ is approximated by
\[

$$
\begin{equation*}
z_{c}^{\prime}(t)=Z_{c}(\infty) \delta(t)+k_{2} \exp \left(-b_{0} t\right)+c_{3} t^{3}+c_{2} t^{2}+c_{1} t+c_{0} \tag{16}
\end{equation*}
$$

\]

### 2.2 Model of propagation function

When $s \rightarrow \infty, \theta(s) d \rightarrow \tau s$ where $\tau=\sqrt{l c} d$ is the propagation delay of the line. In modeling the propagation function, we first extract $\tau s$ from $\theta(s) d$ and let $\phi(s)=$ $\theta(s) d-\tau s$. Then $\Gamma(s)=\Lambda(s) \exp (-\tau s)$ with $\Lambda(s)=\exp (-\phi(s))$. In so doing, for any function $x(s), \Gamma(s) x(s)=\Lambda(s) \exp (-\tau s) x(s)$, and correspondingly, the convolution $\gamma(t) * x(t)=\lambda(t) * x(t-\tau)$. Now we are going to approximate $\lambda(t)$.

As in the case of Sec.2.1, we first compute $\Lambda(0)=\exp (-\sqrt{r g} d), \Lambda(\infty)=\exp \left(-\frac{1}{2} \tau\left(\frac{\tau}{\tau}+\right.\right.$ $\left.\frac{g}{c}\right)$ ) and $\Lambda^{(-1)}=\frac{1}{8} \Lambda(x) \tau\left(\frac{r}{l}-\frac{g}{c}\right)^{2}$. Let the principal part of $\Lambda(s)$ be $\Lambda_{p}(s)=\left(d_{1} s+\right.$ $\left.d_{0}\right) /\left(s+\epsilon_{0}\right)$. then $d_{1}=\Lambda(\infty), \epsilon_{0}=\Lambda^{(-1)} /(\Lambda(0)-\Lambda(\infty))$ and $d_{0}=\Lambda(0) e_{0}$. In the time domain. $\lambda_{p}(t)=d_{1} \delta(t)+k_{\lambda} \epsilon x p\left(-\epsilon_{0} t\right)=\Lambda(\infty) \delta(t)+\Lambda^{(-1)} \exp \left(-e_{0} t\right)$ The remainder $\lambda_{r}(t)=\lambda(t)-\lambda_{p}(t)$ is approximated by $f_{3} t^{3}+f_{2} t^{2}+f_{1} t+f_{0}$ in the same way as in the approximation of the remainder of the characteristic impedance. Then, $\lambda(t) \approx \lambda^{\prime}(t)$ where

$$
\begin{equation*}
\lambda^{\prime}(t)=d_{1} \delta(t)+k_{\lambda} \exp \left(-\epsilon_{0} t\right)+f_{3} t^{3}+f_{2} t^{2}+f_{1} t+f_{0} \tag{17}
\end{equation*}
$$

## Example 1.

Let the frequency $s$ be scaled by a factor $\omega_{0}$ such that $s^{\prime}=s / \omega_{0}$, and let $a=r / \omega_{0} l$, $b=g / \omega_{0}$ and $\tau^{\prime}=\tau \omega_{0}$, then $Z_{c}=Z_{c}(\infty) \sqrt{\left(s^{\prime}+a\right) /\left(s^{\prime}+b\right)}$ and $\Lambda=\exp \left(-\tau^{\prime} \sqrt{\left(s^{\prime}+a\right)\left(s^{\prime}+b\right)}+\right.$ $\left.\tau^{\prime} s^{\prime}\right)$. The accuracy of the approximation of $Z_{c}$ and $\Lambda$ by their principal parts depends on the ratio $a / b$. The near the ratio to 1 , the better the approximation. We show the remainders $z_{c r}(t)$ and $\lambda_{r}(t)$ in Fig.2a and Fig.2b with the ratio $a / b=10^{4}$. The solid lines and dashed lines correspond to the original function and their cubic polynomial fitting. respectively. It can be seen that the fitting is indeed very good.

Remarks.

1. In the case that $g=0$, we can model $y_{c}(t) \equiv z_{c}^{-1}(t)$ instead of modeling $z_{c}(t)$.
2. It can be seen from Fig. 2 that the curve of the remainder is dissymmetric w.r.t. the mid point. A quadratic polynomial can not fit such a curve well, so that we choose a cubic polynomial.
3. The first step of the modeling process to extract a principal part is very essential to the second step. If $z_{c}(t)$ or $\lambda(t)$ is obtained by directly taking the inverse FFT from $Z_{c}(s)$ or $\Lambda(s)$, its curve will have severe ripples. No simple polynomial or piecewise polynomial fitting can be done to approximate such a function well. Also, the numerical computation of a convolution with such a function will be very time consuming and special care must be taken to avoid excessive errors.

### 2.3 Time domain model of the characteristic 2-port

Now we consider the time domain model of the characteristic 2-port. Assuming that the prior simulation time is $t_{n}$ and the current one is $t_{n+1}$ with a step size $h=t_{n+1}-t_{n}$. We consider the the convolution $z_{c}\left(t_{n+1}\right) * i_{1}\left(t_{n+1}\right)$, which is approximated by

$$
\begin{align*}
& z_{c}^{\prime}\left(t_{n+1}\right) * i_{1}\left(t_{n+1}\right)=\int_{0}^{t_{n+1}} z_{c}^{\prime}\left(t_{n+1}-x\right) i_{1}(x) d x= \\
& \int_{0}^{t_{n}} z_{c}^{\prime}\left(t_{n+1}-x\right) i_{1}(x) d x+\int_{t_{n}}^{t_{n+1}} z_{c}^{\prime}\left(t_{n+1}-x\right) i_{1}(x) d x \tag{18}
\end{align*}
$$

The first term of the above expression can be computed recursively as will be described in Sec.4. Let it be expressed as $I_{1 a}\left(t_{n+1}\right)$. When $h$ is small enough, the second term can be approximated by using a trapezoidal formula

$$
\begin{equation*}
\int_{t_{n}}^{t_{n+1}} i_{c}^{\prime}\left(t_{n+1}-x\right) i_{1}(x) d x \approx I_{1 b}\left(t_{n+1}\right)+R(h) i_{1}\left(t_{n+1}\right) \tag{19}
\end{equation*}
$$

where $R(h)=a_{1}+0.5 h\left(k_{z}+c_{0}\right)$ and $I_{1 b}\left(t_{n+1}\right)=0.5 h\left(k_{z} \exp \left(b_{0} h\right)+c_{3} h^{3}+c_{2} h^{2}+c_{1} h+\right.$ $\left.c_{0}\right) i_{1}\left(t_{n}\right)$, which can also be computed. Thus, the convolution $z_{c}\left(t_{n+1}\right) * i_{1}\left(t_{n+1}\right)$ can be expressed as $z(h) i_{1}\left(t_{n+1}\right)+e_{21}\left(t_{n+1}\right)$, where $e_{21}\left(t_{n+1}\right)=I_{1 a}\left(t_{n+1}\right)+I_{1 b}\left(t_{n+1}\right)$, and can be modeled by a resistance $R(h)$ and a voltage source $e_{21}\left(t_{n+1}\right)$ connected in series.

Let $y_{1}(t)$ be $2 r_{1}(t)-u_{1}(t)$. The convolution $\gamma\left(t_{n+1}\right) * y_{1}\left(t_{n+1}\right)=\lambda\left(t_{n+1}\right) * y_{1}\left(t_{n+1}-\tau\right)$. For $t_{n+1} \leq \tau$. it is zero. In the general case. it is easily computed and can be regarded
as a known value. Thus, this convolution can be modeled by an independent voltage source $\epsilon_{\mathrm{w} 1}$.

Therefore, the time domain model of the characteristic 2-port at time $t_{n+1}$ is a resistive two port shown in Fig.3a, where $E_{1, n+1}=e_{x 1}+e_{w 1}$ and $E_{2, n+1}=e_{x 2}+e_{\omega 2}$. This modeled is equivalent to the model shown in Fig.3b, where $J_{1, n+1}=E_{1, n+1} / R(h)$ and $J_{2, n+1}=E_{2, n+1} / R(h)$.

## 3 Model of coupled lines

In the case of a transmission line system with $n$ coupled lines, let $r, l, c$ and $g$ be the resistance, inductance, capacitance and conductance matrix, respectively, and let $V(x, s)$ and $I(x, s)$ be the voltage and current vector, respectively, then the telegrapher's equations have the same forms as those of Eqs.(1) and (2). Now we consider tow kinds of its model.

### 3.1 Characteristic 2n-port model

Based on the above, the characteristic $2 n$-port of $n$ coupled lines can also be described by. Eqs.( 7 ) and ( 8 ), and Eqs.(4) and (5) are also valid. In this case, the characteristic impedance matrix is

$$
\begin{equation*}
Z_{c}(s)=Y^{-1}(s) \sqrt{Y(s) Z(s)} \tag{20}
\end{equation*}
$$

and the propagation function matrix is

$$
\begin{equation*}
\Gamma(s)=\exp (-\sqrt{Y(s) Z(s)} d) \tag{21}
\end{equation*}
$$

Let $\mathrm{Y}^{\prime}(s) Z(s)=Q(s) \Theta^{2}(s) Q^{-1}(s)$, where $\Theta^{2}(s)$ is the diagonal eigen value matrix of $Y(s) Z(s)$ and $T(s)$ its corresponding eigenvector matrix. Then,

$$
\begin{equation*}
\Gamma(s)=Q(s) \exp (-\Theta(s) d) Q^{-1}(s) \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
Z_{\mathrm{c}}(s)=Y^{-1}(s) Q(s) \Theta(s) Q^{-1}(s) \tag{23}
\end{equation*}
$$

$Z_{f}(0)$ can be computed by using Eq.(23). Let

$$
\begin{equation*}
c l=Q(\infty) T^{2} Q^{-1}(\infty) \tag{24}
\end{equation*}
$$

where $T^{2}$ and $Q(\infty)$ are the eigenvalue and eigenvector matrix of the matrix $c l$, respectively. Then, from the above equations, we have

$$
\begin{equation*}
Z_{c}(\infty)=g^{-1} Q(\infty) T Q^{-1}(\infty) \tag{25}
\end{equation*}
$$

The residue matrix $Z_{c}^{(-1)}$ cannot be computed analytically in the general case. To compute an element $Z_{c i j}^{(-1)}$ of $Z_{c}^{(-1)}$, we select two high frequencies $\omega_{1}$ and $\omega_{2}$, and find $v_{1}=Z_{\text {cij }}\left(\omega_{1}\right)$ and $v_{2}=Z_{c i j}\left(\omega_{2}\right)$. Let $u_{k}=\operatorname{imag}\left(v_{k}\right)$ for $k=1,2$. Then,

$$
\begin{equation*}
Z_{\mathrm{cij}}^{(-1)}=\frac{u_{1} \omega_{1}^{3}-u_{2} \omega_{2}^{3}}{\omega_{2}^{2}-\omega_{1}^{2}} \tag{26}
\end{equation*}
$$

We can use the formulas given in Sec.(2.1) to form the principal part of each element of the characteristic impedance matrix, and then to approximate its remainder by using a piecewise cubic fitting as will be described in the next section.

Let $Q(s)=\left[Q_{i j}(s)\right], Q^{-1}(s)=\left[R_{i j}(s)\right]$ and $\Theta(s)=\operatorname{diag}\left[\Theta_{i}(s)\right]$, then

$$
\begin{equation*}
\Gamma_{i j}(s)=\sum_{k=1}^{n} S_{i j k} \exp \left(-\Theta_{k}(s) d\right) \tag{27}
\end{equation*}
$$

where $S_{i j k}=Q_{i k} R_{k j}$. Let $T=\operatorname{diag}\left[\tau_{k} / d\right]$ and $\Phi_{k}(s)=\Theta_{k}-s T_{k}$, then

$$
\begin{equation*}
\Gamma_{i j}(s)=\sum_{k=1}^{n}\left[\exp \left(-\tau_{k} s\right) S_{i j k} \in x p\left(-\Phi_{k}(s) d\right)\right] \tag{28}
\end{equation*}
$$

In the general case, the $n$ components of $\Gamma_{i j}(s)$ have different ideal delay factor $\exp \left(-\tau_{k} s\right)$ and each of them need to be modeled individually.

From what was mentioned above, there are $n^{2}$ elements in the model of $z_{c}(t)$ and $n^{3}$ elements in the model of $\gamma(t)$, and there are $2\left(n^{2}+n^{3}\right)$ convolutions to take with these two functions.

The time domain characteristic 2-port model with two coupled lines is shown in Fig. 4 for illustration. In the figure, $R_{11(h)}$ and the CCVS $R_{12} I_{12, n+1}$ come from the convolution $z_{11} * i_{11}+z_{12} * i_{12}$, and the independent voltage source $E_{11, n+1}$ comes from the abore convolution and the convolution $\gamma_{11} *\left(2 v_{21}-w_{21}\right)+\gamma_{12} *\left(2 v_{22}-w_{22}\right)$. Other elements come from similar convolutions.

### 3.2 Model with decoupling transformation

Another way to model a system of coupled transmission lines is to use decoupling transformations and to model each decoupled single line by its characteristic 2-port
[10]. Here we follow the same idea as in [10] with a slight variation.
Let $A=Z^{1 / 2} Y^{\prime} Z^{1 / 2}=W \cdot \Theta^{2} W^{-1}$ where $\Theta^{2}=\operatorname{diag}\left[\theta_{k}^{2}\right]$ is the diagonal eigenvalue matrix of $A$, and $W$ is the corresponding eigenvector matrix. Let

$$
\begin{gather*}
X(s)=Z^{1 / 2} W \operatorname{diag}\left[\sqrt{1 / z_{0 k} \theta_{k}}\right]  \tag{29}\\
P(s)=\operatorname{diag}\left[\sqrt{1 / z_{0 k} \theta_{k}}\right] W^{-1} Z^{1 / 2} \tag{30}
\end{gather*}
$$

$V(x, s)=X(s) E(x, s)$ and $J(x, s)=P(s) I(x, s)$. Then, we have the decoupled system

$$
\begin{align*}
& \frac{d E(x, s)}{d x}=-\tilde{Z}(s) J(x, s)  \tag{31}\\
& \frac{d J(x, s)}{d x}=-\tilde{Y}(s) E(x, s) \tag{32}
\end{align*}
$$

where $\dot{Z}(s)=\operatorname{diag}\left(z_{0 k} \theta_{k}\right)$ and $\tilde{Y}(s)=\operatorname{diag}\left(\theta_{k} / z_{0 k}\right)$. The characteristic impedance and the propagation function of the $k$-th decoupled line are $z_{0 k}$ and $\exp \left(-\theta_{k} d\right)$, respectively. Note that $z_{0 k}$ can be arbitrarily chosen. For simplicity, we just choose $z_{0 k}=1$ for each k. ${ }^{2}$

Thus. by using the decoupling transformations, the modeling of a coupled line system can be done by modeling its decoupling transformation matrices and each of the decoupled single lines. As the characteristic impedance of each decoupled line is a constant. we only need to model its propagation function.

In the formation of the principal parts of the transformation matrices, $X(0)$ and $P(0)$ can be computed by using Eqs.(29) and (30). In order to compute their values at $s=x$. let $l^{1 / 2} c l^{1 / 2} \equiv B=W_{\infty} T^{2} W_{\infty}^{-1} \cdot{ }^{3}$ Then,

$$
\begin{equation*}
X(\infty)=l^{1 / 2} W_{\infty} \operatorname{diag}\left[\sqrt{d / z_{0 k} \tau_{k}}\right] \tag{33}
\end{equation*}
$$

and

$$
\begin{equation*}
P(\infty)=\operatorname{diag}\left[\sqrt{d / z_{0 k} \tau_{k}}\right] W_{\infty}^{-1} l^{1 / 2} \tag{34}
\end{equation*}
$$

The residues $X^{(-1)}$ and $P^{(-1)}$ can be computed numerically as described in Sec.3.1.

[^1]There are two cases encountered in the formation of the principal parts of the transformation matrices.

## Case. 1

In the formation of the principal part $F_{p}=\left(a_{1} s+a_{0}\right) /\left(s+b_{0}\right), b_{0}=F^{(-1)} /(F(0)-$ $F(\infty)$ ). In order that the model be stable, $F^{(-1)}$ and $F(0)-F(\infty)$ must have the same sign. When this condition is violated, we use a technique called shift at zero frequency (SZF) to let the principal part be $\left(a_{1} s+a_{0}\right) /\left(s+b_{0}\right)-c_{0} /\left(d_{1} s+1\right)$. In the additional term, $c_{0}$ is so chosen that $F^{(-1)}$ and $F(0)+c_{0}-F_{\infty}$ has the same sign and $b_{0}=F^{(-1)} /\left(F(0)+c_{0}-F(\infty)\right)$ is positive. $d_{1}$ is set as large as possible so that the second term has no real effect on the frequency response for $f \geq f_{0}$ and $F^{(-1)}$ and $F(\infty)$ remain unchanged. The inverse Laplace transform of the additional term is $-c_{0} / d_{1} \epsilon x p\left(-t / d_{1}\right)$. As $d_{1}$ is very large, practically speaking, it can be neglected in the time domain model.

Case 2.
This case may happen in dealing with the numerical computation of residues.
Let $\mu_{\text {max }}=2 \pi f_{\text {max }}$ and $f\left(j u_{\max }\right)=u+j v$ where $u$ and $v$ are real values. Then, $F\left(x^{-}\right) \approx u . F^{(-1)} \approx-v * \omega_{\max }$, and the coefficient $b_{0} \approx-v * \omega_{\max } /(F(0)-u)$. In order that the principal part approximates $F(s)$ well in high frequencies, $\left|b_{0}\right|$ should be much smaller than $u_{\text {max }}$, i.e. $|r| \ll|F(0)-F(\infty)|$ is needed. We have found that this is not always the case. When the opposite case takes place, a technique called shift at infinity frequency (SIF) similar to SZF can be used. In SIF, we modify $F_{p}(s)$ by subtraction of $Y(s)=q_{1} s /\left(s+q_{0}\right)$. When $s \rightarrow \propto, Y(s) \rightarrow q_{1} . q_{1}$ is so chosen that $|v| \ll\left|F(0)-F(\infty)-q_{1}\right|$ is satisfied. $q_{0}$ is so chosen that $\left|q_{0}\right| \ll f_{0}$ and $L^{-1} y(s) \approx q_{1} \delta(t)$.

A time domain model of a two coupled line system made from the decoupling transformation is shown in Fig. 5 for illustration. In the figure, the subscript " $n+1$ " used to specify the time is omitted for brevity. The VCVS $S_{11}$ and the independent voltage source $T_{11}$ come from the convolution $x_{11} * \epsilon_{11}+x_{12} * e_{12}$. The CCCS $J_{11}$ and the independent current source $U_{11}$ come from the convolution $p_{11} * j_{11}+p_{12} * j_{12}$. $R$ is the characteristic impedance of the decoupled lines, and $W_{11}$ comes from the convolution $\hat{h}_{11} *\left(2 \epsilon_{21}-u_{21}\right)$. Other elements are from similar convolutions.

Now we compare the models without and with decoupling transformations, which are called model 1 and model 2 for brevity. In model 2, there are $n$ elements in the model of propagation functions, and $2 n^{2}$ elements in the model of transformation matrices. The number of convolution related to these elements are $2 n+4 n^{2}$. They are about a factor of $2 / n$ w.r.t. those of model 1 . On the other hand, there are $6 n$ unknowns in the circuit model with model 2 ( 4 n voltages and 2 n currents), compared with 4 n unknowns ( 2 n voltages and 2 n currents) in model 1 . For $n \geq 3$, the model with decoupling transformations is more efficient.

## 4 Piecewise cubic fitting

In the formation of the coupled line model, we use piecewise cubic fitting to approximate a remainder.

Let $f(t)$ be the function to be approximated, which is defined in the time interval $\left[0, t_{\text {mar }}\right]$. Let the interval be divided inton subintervals with breaking points $0 \equiv t_{0}<$ $t_{1}<t_{2}, \ldots<t_{n-1}<t_{n} \equiv t_{\text {mar }}$. In the m-th subinterval $I_{m}=\left[t_{m-1}, t_{m}\right]$, suppose that $f(t)$ is approximated by $f_{m}(t)=a_{3}^{m}\left(t-t_{m-1}\right)+a_{2}^{m}\left(t-t_{m-1}\right)^{2}+a_{1}^{m}\left(t-t_{m-1}\right)+a_{0}^{m}$. For $m>1$, we let $a_{0}^{m}=f_{m-1}\left(t_{m-1}\right)$ so that the piecewise cubic polynomials are continuous at the breaking points. Let $t_{i}^{m}$ be the j -th sampling point in the m -th subinterval. Suppose that there are $s_{m}$ sampling points. Then, the least square fitting of $f_{m}(t)$ to $f(t)$ in the interval $I_{m}$ is equivalent to finding the least square solution of the linear set of equations

$$
\begin{equation*}
H^{m} X^{m}=J^{m} \tag{35}
\end{equation*}
$$

where $X^{1}=\left[a_{0}^{1} \cdot a_{1}^{1}, a_{2}^{1}, a_{3}^{1}\right]^{t}$ and $X^{m}=\left[a_{1}^{m}, a_{2}^{m}, a_{3}^{m}\right]^{t}$ for $m>1 . H^{1}$ is an $s_{1} \times 4$ matrix with its element $h_{i j}^{1}=\left(t_{i}^{1}\right)^{j-1}$ and $J^{1}$ is an $s_{1}$ vector with its element $J_{i}^{1}=f\left(t_{i}^{1}\right)$. For $m>1, H^{m}$ is an $s^{m} \times 3$ matrix with its element $h_{i j}^{m}=\left(t_{i}^{m}\right)^{j}$ and the i-th element of $J^{m}$ is $J_{i}^{m}=f\left(t_{1}^{m}\right)-a_{0}^{m}$. By using the Householder transformation to transform Eq.(35) to the following form

$$
\left[\begin{array}{c}
G^{m}  \tag{36}\\
0
\end{array}\right] X^{m}=\left[\begin{array}{l}
B^{m} \\
C^{m}
\end{array}\right]
$$

where $G^{m}$ is an upper triangular matrix, then $X^{m}$ can be found by solving the equations $G^{m} X^{m}=B^{m}$.

The piecewise cubic fitting is done step by step from the first interval to the last one, and the breaking points $0 \equiv t_{0}<t_{1}<t_{2}, \ldots<t_{n-1}<t_{n} \equiv t_{\text {max }}$ are determined in the process one after another. We first divide the interval $\left[0, t_{\text {max }}\right]$ into $p$ monotonic subintervals $I_{1}, I_{2}, \ldots, I_{p}$ with $I_{k}$ defined by the terminal points [ $d_{k-1}, d_{k}$ ], where $d_{0}=t_{0}$ and $d_{p}=t_{\text {max }}$. As a cubic polynomial may have at most two extremes, the initial position of the first breaking point $t_{1}$ is set in $I_{2} \operatorname{such}$ that $t_{1}=\min \left(1.2 d_{1}, d_{2}\right)$. Suppose that the $k$-th breaking point $t_{k}$ is in $I_{j}$, then the initial position of the $k+1$-th breaking point is set at $t_{k+1}=\min \left(t_{k}+1.2 *\left(d_{j}-t_{k}\right), d_{j+1}, t_{\max }\right)$. After the initial position of the breaking point of a new subinterval is set, a least square fitting is done, and the root mean square of the relative error $r m s$ is computed. If $r m s$ is within certain limit (e.g.: $4 \%$ to $5.5 \%$ ), then the new breaking point is set. A farther or nearer breaking point is tried depending on the case whether rms is too small or too big. In either case. the change of the length of the subinterval is limited by a factor of 2 or $1 / 2$, and the number of iterations is limited by 4 . This process continues until the boundary of the interval is reached.

After the piecewise cubic fitting is done, the piecewise cubic polynomial is transformed into the form of $\sum_{k=0}^{m} g_{k}\left(t-t_{k}\right) 1\left(t-t_{k}\right)$ where $g_{k}\left(t-t_{k}\right)=a_{3}^{k}\left(t-t_{k}\right)^{3}+a_{2}^{k}(t-$ $\left.t^{k}\right)^{2}+a_{1}^{k}\left(t-t_{k}\right)+a_{0}^{k} .1\left(t-t_{k}\right)$ is a unit step function starting at $t_{k}$ and $a_{0}^{k}=0$ for $k>0$.

## Example 2.

In Fig.6-8, we show an example of the remainders of the propagation functions and the elements of the decoupling matrices $P$ and $X$ of a two coupled line system. The solid and dashed lines correspond to the original functions and their piecewise cubic approximations. It can be seen that the approximations are quite good.

## 5 Recursive convolution

Each element of our model consists of an impulse, an exponential function, and a number of cubic polynomials starting at different time points. When such a function
is convolved with another function $x(t)$, the convolution can be done recursively from time to time. The recursive convolution formula with an exponential function has been given in [ 7 ], and the recursive formulas with polynomials up to the order of 2 have been given in [13]. We now give the recursive formulas with a cubic polynomial for reference.

Let $f_{k}(t, T)=(t-T)^{k} 1(t-T)$ and $I_{k}(t, x)=f_{k}(t, T) * x(t)=\int_{0}^{t}(t-T-\tau)^{k} 1(t-$ $T-\tau) x(\tau) d \tau$. Then, for the time $t+h>T$, when trapezoidal formula is used to do integration, we have the recursive formulas for the convolution from $k=0$ to 3 as follows:

$$
\begin{aligned}
& I_{0}(t+h, x)=I_{0}(t, x)+0.5 h(x(t-T)+x(t+h-T)), \\
& I_{1}(t+h, x)=I_{1}(t, x)+h I_{0}(t, x)+0.5 h^{2} x(t-T), \\
& I_{2}(t+h, x)=I_{2}(t, x)+2 h I_{1}(t, x)+h^{2} I_{0}(t, x)+0.5 h^{3} x(t-T), \\
& I_{3}(t+h, x)=I_{3}(t, x)+3 h I_{2}(t, x)+3 h^{2} I_{1}(t, x)+h^{3} I_{0}(t, x)+0.5 h^{4} x(t-T) .
\end{aligned}
$$

Note that when $T>0$, these convolutions have no relation with the current value of $x$. i.e., $x(t+h)$ : and when $T=0$, only $I_{0}(t+h, x)$ has a term related to $x(t+h)$.

## 6 Examples and conclusion

We present two examples to show the simulation results with our model. The first example is a single line circuit shown in Fig.9, with the line voltages shown in Fig.10. The second example is a two-coupled line circuit shown in Fig.11, with the line voltages shown in Fig.12. The solid lines correspond to the simulation results by using our model and the recursive convolution, while the dashed lines (with an extension ".ff") correspond to those obtained by directly using FFT and exact expressions of the characteristics of the lines. It can be seen that these two results match very well, which shows the accuracy of our model.

We also test the case when the remainders of the characteristics are not approximated by piecewise cubic polynomial so that no recursive convolution can be used. The CPC゙ time is almost 100 times more than using our model with recursive convolution, which shows the efficiency of our model.

We have presented a new time domain model for single and coupled transmission lines. The model of each element consists of an impulse function, an exponential func-
tion, and a piecewise cubic polynomial. The model is stable. As very good approximation can be obtained in the piecewise cubic polynomial fitting for the remainders, the model can be very accurate. In fact, as the principal part approximates the characteristic well in a wide frequency region, the remainder is relatively small and only a few pieces of piecewise cubic polynomials can lead to very good accuracy. Therefore, the model is both simple and accurate. As the model is compatible with recursive convolution, it is efficient in the use of time domain simulation. Compared with the model formed by using least square fitting in frequency domain only, our model is easier to form and computationally inexpensive. Our future work is to extend our model to nonuniform and frequency-dependent lines.

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## Captions of Figures 2, 6,7,8, 10 and 12

Fig. 2 Remainders of characteristic impedance and propagation function of a single line

Fig. 6 Remainders of propagation function of two coupled lines

Fig. 7 Remainders of decoupling transformation matrix X of two coupled lines

Fig. 8 Remainders of decoupling transformation matrix $P$ of two coupled lines

Fig. 10 Simulation result of a single line circuit
Fig. 12 Simulation result of two coupled line circuit


Fig.2a








Fig.6a


Fig.6b


Fig.7a



Fig.7e


Fig.7d



Fig.8b



Fig.8d



$\mathrm{Y} \times 10^{-3}$
Fig.10a


Fig.10b



Fig.12b


Fig.12c


Fig.12d



[^0]:    ${ }^{1}$ We use the superscript ( -1 ) to denote the residue of a function and the superscript -1 without paratheses to denote an inverse.

[^1]:    ${ }^{2}$ At frequencies 0 and $\propto$, our formulas are the same as given in [10], but at a frequency $s=j \omega$, they are different. The formulas given in [10] is based on the assumption that matrix $A$ is real symmetric, so that matrix $W$ is orthonormal. This is not the case when $s=j u$ for finite nonzero $\omega$.
    ${ }^{3}$ The matrix $T$ here is the same as in Eq.(24)

