# LOGARITHMIC POTENTIALS AND PLANAR BROWNIAN MOTION 

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In this paper we continue our discussion of the connection between potential theory and Brownian motion begun in "Classical Potential Theory and Brownian Motion" that also appears in this Symposium volume. Throughout this paper, we will be dealing with a two dimensional Brownian motion process. We will continue numbering the sections from where we left off in the previous paper.

## 8. Planar Brownian motion

In Section 5, we saw that for a Brownian motion process in $n \geqq 3$ dimensions, $P_{x}\left(\lim _{t \rightarrow \infty}\left|X_{t}\right|=\infty\right)=1$ for all $x$. In sharp contrast to this situation, a planar Brownian motion is certain to hit any nonpolar set.

Theorem 8.1. Let $B$ be a Borel set. Then $P_{x}\left(V_{B}<\infty\right)$ is either identically 1 or identically 0 .
Proof. A simple computation shows that for any $x \in R^{2}, \int_{0}^{t} p(s, x) d s \uparrow \infty$ as $t \uparrow \infty$. Thus, for any nonnegative function $f$ having nonzero integral,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \int_{0}^{t} P^{s} f(x) d s=\infty \tag{8.1}
\end{equation*}
$$

Let $\varphi(x)=P_{x}\left(V_{B}<\infty\right)$. Then for any $h>0$,

$$
\begin{equation*}
0 \leqq \int_{0}^{t} P^{s}\left(\varphi-P^{h} \varphi\right) d s=\int_{0}^{h} P^{s} \varphi d s-\int_{t}^{t+h} P^{s} \varphi d s \leqq 2 h \tag{8.2}
\end{equation*}
$$

Letting $t \uparrow \infty$, we see that

$$
\begin{equation*}
0 \leqq \int_{0}^{\infty} P^{s}\left(\varphi-P^{h} \varphi\right) d s \leqq 2 h \tag{8.3}
\end{equation*}
$$

But then it must be that $\varphi=P^{h} \varphi$ a.e. Since $P^{t} \varphi \uparrow \varphi$ as $t \downarrow 0$ and $P^{t}\left(P^{h} \varphi\right) \uparrow P^{h} \varphi$ as $t \downarrow 0$, it follows that $\varphi(x)=P^{h} \varphi(x)$ for all $x$. Using Proposition 2.3, we see that $\varphi(x) \equiv \alpha$ for some constant $\alpha$. Now

$$
\begin{equation*}
P_{x}\left(t<V_{B}<\infty\right)=\int_{R^{2}} q_{B}(t, x, y) \varphi(y) d y=\alpha P_{x}\left(V_{B}>t\right) \tag{8.4}
\end{equation*}
$$

[^0]Letting $t \uparrow \infty$, we see that $\alpha P_{x}\left(V_{B}=\infty\right)=0$. Thus, either $P_{x}\left(V_{B}=\infty\right) \equiv 0$ or $\alpha=0$. In the first case $\varphi(x) \equiv 1$, while in the second case $\varphi(x) \equiv 0$. This establishes the proposition.

The difference between planar Brownian motion and Brownian motion in $n \geqq 3$ dimensions has its analytical counterpart in potential theory. We will now show that the potentials associated with planar Brownian motion are logarithmic potentials.

Let 1 denote the point $(1,0)$ and let $a^{\lambda}(x)=g^{\lambda}(1)-g^{\lambda}(x)$. Using (2.30), we see that for $x \neq y$,

$$
\begin{equation*}
a^{\lambda}(y-x)=\int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)-g_{B}^{\lambda}(x, y)+L_{B}^{\lambda}(x), \tag{8.5}
\end{equation*}
$$

where

$$
\begin{equation*}
L_{B}^{\lambda}(x)=g^{\lambda}(1)\left[1-E_{x}\left(\exp \left\{-\lambda V_{B}\right\}\right)\right] \tag{8.6}
\end{equation*}
$$

Now

$$
\begin{equation*}
a^{\lambda}(x)=\int_{0}^{\infty} e^{-\lambda t} \cdot[p(t, \mathbf{l})-p(t, x)] d t \tag{8.7}
\end{equation*}
$$

If $|x| \geqq 1$, then $p(t, 1)-p(t, x) \geqq 0$; so for $|x| \geqq 1, a^{\lambda}(x)$ is increasing. On the other hand, for $|x|<1, p(t, 1)-p(t, x)<0$ so $-a^{\lambda}(x)$ is increasing. In either case

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} a^{\lambda}(x)=\int_{0}^{\infty}[p(t, 1)-p(t, x)] d t=\frac{1}{\pi} \log |x| \tag{8.8}
\end{equation*}
$$

and the convergence is uniform on any compact set not containing 0 . For simplicity, we set $a(x)=(1 / \pi) \log |x|$.

Our principle result in this section will be to establish the following theorem.
Theorem 8.2. Let $B$ be a nonpolar set. Then $g_{B}(x, y)<\infty$ for $x \neq y$ and $\lim _{\lambda \downarrow 0} L_{B}^{\lambda}(x)=L_{B}(x)$ exists and is finite for all $x$. Moreover, for $x \neq y$,

$$
\begin{equation*}
a(y-x)-\int_{\bar{B}} \Pi_{B}(x, d z) a(y-z)=-g_{B}(x, y)+L_{B}(x) . \tag{8.9}
\end{equation*}
$$

Before getting on with the proof, we observe first if $E_{x} V_{B}<\infty$ for all $x$ then $L_{B}(x) \equiv 0$. Indeed,

$$
\begin{equation*}
L_{B}^{\lambda}(x)=\lambda g^{\lambda}(1)\left[\frac{1-E_{x}\left(\exp \left\{-\lambda V_{B}\right\}\right)}{\lambda}\right] \tag{8.10}
\end{equation*}
$$

so if $E_{x} V_{B}<\infty$, then the expression in the brackets converges to $E_{x} V_{B}<\infty$, while $\lambda g^{\lambda}(1) \rightarrow 0$ as $\lambda \rightarrow 0$. In particular by Proposition $2.2, L_{B}(x) \equiv 0$ whenever $B^{c}$ is relatively compact.

The proof of Theorem 8.2 is long and will be divided into several lemmas.
Lemma 8.1. Suppose B is relatively compact. Then

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{\bar{B}} \Pi_{B}(x, d z) a(y-z) . \tag{8.11}
\end{equation*}
$$

Proof. If $B$ is polar, there is nothing to prove since both $\Pi_{B}^{\lambda}$ and $\Pi_{B}$ are the zero measure, so suppose $B$ is nonpolar. If $y \notin \bar{B}$, then as $\lambda \downarrow 0, a^{\lambda}(y-z)$ converges to $a(y-z)$ uniformly in $z \in \bar{B}$, and thus (8.11) holds in this case. Suppose $y \in \bar{B}$. Let $D_{\varepsilon}$ be the open disk of center $y$ and radius $\varepsilon<1$. We can write

$$
\begin{align*}
& \int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)  \tag{8.12}\\
& \quad=\int_{\bar{B} \cap D_{\bar{E}}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)+\int_{\bar{B} \cap D_{\varepsilon}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z) .
\end{align*}
$$

Since

$$
\begin{equation*}
\int_{\bar{B} \cap D \in} \prod_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=E_{x}\left[\exp \left\{-\lambda V_{B}\right\} a^{\lambda}\left(y-X_{V_{B}}\right) 1_{D \varepsilon}\left(X_{V_{B}}\right)\right] \tag{8.13}
\end{equation*}
$$

and $a^{\lambda}(y-z)$ converges to $a(y-z)$ uniformly for $z \in \bar{B} \cap D_{\varepsilon}^{c}$, we see that

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \int_{\bar{B}_{\cap} \cap \bar{\varepsilon}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=E_{x}\left[a\left(y-X_{V_{B}}\right) 1_{D \varepsilon}\left(X_{V_{B}}\right)\right] . \tag{8.14}
\end{equation*}
$$

On the other hand,

$$
\begin{align*}
-\int_{\bar{B} \cap D_{\varepsilon}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z) &  \tag{8.15}\\
& =E_{x}\left[\exp \left\{-\lambda V_{B}\right\}\left[-a^{\lambda}\left(y-X_{V_{B}}\right)\right] 1_{D_{\varepsilon}}\left(X_{V_{B}}\right) .\right.
\end{align*}
$$

Since

$$
\begin{equation*}
-\exp \left\{-\lambda V_{B}\right\} a^{\lambda}\left(y-X_{V_{B}}\right) 1_{D_{\varepsilon}}\left(X_{V_{B}}\right) \uparrow-a\left(y-X_{V_{B}}\right) 1_{D_{\varepsilon}}\left(X_{V_{B}}\right), \tag{8.16}
\end{equation*}
$$

monotone convergence shows that

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \int_{\bar{B} \cap D_{\varepsilon}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=E_{x}\left[a\left(y-X_{V_{B}}\right) 1_{D_{\varepsilon}}\left(X_{V_{B}}\right)\right] \geqq-\infty . \tag{8.17}
\end{equation*}
$$

As

$$
\begin{equation*}
-\infty<E_{x}\left[a\left(y-X_{V_{B}}\right) \mathbf{1}_{D_{\varepsilon}}\left(X_{V_{B}}\right)<+\infty,\right. \tag{8.18}
\end{equation*}
$$

we see from (8.14) and (8.17) that (8.11) holds. This establishes the lemma.
Lemma 8.2. Suppose $B$ is a nonpolar relatively compact set. Let $A$ be a compact set of positive measure such that $A \cap \bar{B}=\varnothing$. Then

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \int_{A} g_{B}^{\lambda}(x, y) d y=E_{x} \int_{0}^{V_{B}} 1_{A}\left(X_{t}\right) d t<\infty \tag{8.19}
\end{equation*}
$$

Proof. For a given $x_{0}$ there is a $t_{0}>1$ such that

$$
\begin{align*}
& P_{x_{0}}\left(X_{s} \in B \text { for some } s \in\left(1, t_{0}\right)\right)  \tag{8.20}\\
& \quad=\int_{R^{2}} p\left(1, y-x_{0}\right) P_{y}\left(V_{B} \leqq t_{0}-1\right) d y>0,
\end{align*}
$$

since otherwise

$$
\begin{equation*}
\int_{R^{n}} p\left(1, y-x_{0}\right) P_{y}\left(V_{B}<\infty\right) d y=P_{x_{0}}\left(V_{B}<\infty\right)=0 . \tag{8.21}
\end{equation*}
$$

Thus, there must be a compact set $F$ of positive measure such that $P_{y}\left(V_{B} \leqq\right.$ $\left.t_{0}-1\right)>0$ for all $y \in F$. As $p(1, x)$ is a strictly positive continuous function, it follows that

$$
\begin{equation*}
\inf _{x \in A} P_{x}\left(V_{B} \leqq t_{0}\right) \geqq \inf _{x \in A} \int_{F} p(1, y-x) P_{y}\left(V_{B} \leqq t_{0}-1\right)=\delta>0 \tag{8.22}
\end{equation*}
$$

Let $I_{j}=\left[j t_{0},(j+1) t_{0}\right)$ and let $C=\left\{t: X_{t} \in A, V_{B}>t\right\}$. Define the index set $\Gamma$ by $j \in \Gamma$ if and only if $I_{j} \cap C \neq \varnothing$, and enumerate $\Gamma$ increasing order by $j_{1}<j_{2}<\cdots$. Define the times $T_{1}<T_{2}<\cdots$ as follows:

$$
\begin{equation*}
T_{1}=\inf \{t: t \in C\}(=\infty \text { if there is no such } t) \tag{8.23}
\end{equation*}
$$

and

$$
\begin{equation*}
T_{n+1}=\inf \left\{t: t \in C \cap\left[j_{n} t_{0}, \infty\right\}(=\infty \text { if there is no such } t)\right. \tag{8.24}
\end{equation*}
$$

Let $N \leqq \infty$ denote the number of indices in $\Gamma$. Then

$$
\begin{align*}
P_{x}(N>n, N \leqq n+2) & =P_{x}\left(T_{n}<\infty, T_{n+2}=\infty\right)  \tag{8.25}\\
& \geqq P_{x}\left(T_{n}<\infty, V_{B} \leqq T_{n}+t_{0}\right) \\
& =\int_{A} P_{x}\left(T_{n}<\infty, X_{T_{n}} \in d z\right) P_{z}\left(V_{B} \leqq t_{0}\right) \\
& \geqq \delta P_{x}\left(T_{n}<\infty\right)=\delta P_{x}(N>n),
\end{align*}
$$

so $P_{x}(N>n+2) \leqq(1-\delta) P_{x}(N>n)$. Thus, $E_{x} N<\infty$, and hence

$$
\begin{equation*}
E_{x} \int_{0}^{V_{B}} 1_{A}\left(X_{t}\right) d t=E_{x}|C| \leqq E_{x}\left|\bigcup_{j \in \Gamma} I_{j}\right|=t_{0} E_{x} N<\infty, \tag{8.26}
\end{equation*}
$$

as desired. This establishes Lemma 8.2.
We can now prove Theorem 8.2 when $B$ is a relatively compact set.
Lemma 8.3. Suppose $B$ is a nonplanar relatively compact set. Then Theorem 8.2 holds.

Proof. Let $A$ be as in Lemma 8.2. Since

$$
\begin{equation*}
\int_{A} a^{\lambda}(y-x) d y \rightarrow \int_{A} a(y-x) d y \tag{8.27}
\end{equation*}
$$

uniformly in $x$ on compacts, we see that

$$
\begin{equation*}
\lim _{\lambda \downharpoonright 0} \int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) \int_{A} a^{\lambda}(y-z) d y=\int_{\bar{B}} \Pi_{B}(x, d z) \int_{A} a(y-z) d y . \tag{8.28}
\end{equation*}
$$

By Lemma 8.2,

$$
\begin{equation*}
\int_{A} g_{B}^{\lambda}(x, y) d y \uparrow \int_{A} g_{B}(x, y) d y<\infty \tag{8.29}
\end{equation*}
$$

By (8.5)

$$
\begin{align*}
&|A| L_{B}^{\lambda}(x)=\int_{A} a^{\lambda}(y-x) d y-\int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) \int_{A} a^{\lambda}(y-z) d y  \tag{8.30}\\
&+\int_{A} g_{B}^{\lambda}(x, y) d y
\end{align*}
$$

Since the right side has a finite limit as $\lambda \downarrow 0$, we see that $L_{B}^{\lambda}(x)$ must have a finite limit as $\lambda \downarrow 0$. Call this limit function $L_{B}(x)$. By (8.5),

$$
\begin{equation*}
-g_{B}^{\lambda}(x, y)=a^{\lambda}(y-x)-\int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)-L_{B}^{\lambda}(x) \tag{8.31}
\end{equation*}
$$

and thus for $x \neq y$,

$$
\begin{equation*}
\lim _{\lambda \downarrow 0}\left[-g_{B}^{\lambda}(x, y)+\int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)\right]=a(y-x)-L_{B}(x) \tag{8.32}
\end{equation*}
$$

is finite. Now $0 \leqq g_{B}^{\lambda}(x, y) \uparrow g_{B}(x, y) \leqq+\infty$ and by Lemma 8.1

$$
\begin{equation*}
\lim _{\lambda \downharpoonright 0} \int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{\bar{B}} \Pi_{B}(x, d z) a(y-z), \tag{8.33}
\end{equation*}
$$

and

$$
\begin{equation*}
-\infty \leqq \int_{\bar{B}} \Pi_{B}(x, d z) a(y-z)<+\infty . \tag{8.34}
\end{equation*}
$$

Thus for $x \neq y$, we see that $g_{B}(x, y)<+\infty, \int_{\bar{B}} \Pi_{B}(x, d z) a(y-z)>-\infty$, and that (8.9) holds.

To handle the unbounded case we need two additional lemmas.
Lemma 8.4. Suppose $A \subset B$. Then $g_{A}(x, y) \geqq g_{B}(x, y)$ for all $x, y$.
Proof. Since $A \subset B, V_{A} \geqq V_{B}$ and thus for any Borel set $F$,

$$
\begin{equation*}
P_{x}\left(V_{A}>t, X_{t} \in F\right) \geqq P_{x}\left(V_{B}>t, X_{t} \in F\right) . \tag{8.35}
\end{equation*}
$$

Hence (in the notation of Section 2) for a.e. $y$,

$$
\begin{equation*}
q_{A}(t, x, y) \geqq q_{B}(t, x, y) . \tag{8.36}
\end{equation*}
$$

But then

$$
\begin{equation*}
\int_{R^{2}} q_{A}(t-\varepsilon, x, z) p(\varepsilon, y-z) d z \geqq \int_{R^{2}} q_{B}(t-\varepsilon, x, z) p(\varepsilon, y-z) d z, \tag{8.37}
\end{equation*}
$$

and thus letting $\varepsilon \downarrow 0$ (see Section 2), it follows that (8.36) holds for all $y$.

Integrating on $t$, we see that

$$
\begin{equation*}
g_{A}(x, y)=\int_{0}^{\infty} q_{A}(t, x, y) d t \geqq \int_{0}^{\infty} q_{B}(t, x, y) d t=g_{B}(x, y) \tag{8.38}
\end{equation*}
$$

as desired.
Lemma 8.5. Let $B$ be any Borel set. Then for $x \neq y$,

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{\bar{B}} \Pi_{B}(x, d z) a(y-z)>-\infty . \tag{8.39}
\end{equation*}
$$

Proof. If $B$ is polar, there is nothing to prove so suppose $B$ is nonpolar. Suppose $y \in(\bar{B})^{c}$. Let $D_{r}$ be a disk of center 0 and radius $r>1+|y|$. Then $|y-z|>1$ for $z \in D_{r}^{c}$. Hence, by monotone convergence

$$
\begin{gather*}
\int_{\bar{B} \cap D F} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=E_{x}\left[\exp \left\{-\lambda V_{B}\right\} a^{\lambda}\left(y-X_{V_{B}}\right) \mathbf{1}_{D F}\left(X_{V_{B}}\right)\right]  \tag{8.40}\\
\uparrow E_{x}\left[a\left(y-X_{V_{B}}\right) 1_{D F}\left(X_{V_{B}}\right)\right]=\int_{\bar{B} \cap D F} \Pi_{B}(x, d z) a(y-z) \geqq 0 .
\end{gather*}
$$

On the other hand, $a^{\lambda}(y-z) \rightarrow a(y-z)$ uniformly in $z \in \bar{B} \cap D_{r}$ so

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \int_{\bar{B} \cap D_{r}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{\bar{B} \cap D_{r}} \Pi_{B}(x, d z) a(y-z), \tag{8.41}
\end{equation*}
$$

and

$$
\begin{equation*}
-\infty<\int_{\bar{B} \cap D_{r}} \Pi_{B}(x, d z) a(y-z)<\infty \tag{8.42}
\end{equation*}
$$

Thus, (8.39) holds for $y \notin \bar{B}$. Suppose $y \in \bar{B}$ and let $A_{\varepsilon}$ be the disk of center $y$ and radius $\varepsilon<1$. We can write

$$
\begin{align*}
\int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{D_{r} \cap A_{\varepsilon}} & \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)  \tag{8.43}\\
& +\int_{D_{r} \cap A \varepsilon} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z) \\
& +\int_{D F} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z) .
\end{align*}
$$

Choose $r$ so large that $B \cap D_{r}$ is nonpolar and such that $|y-z|>1$ for $z \in D_{r}^{c}$. The second term on the right of (8.43) converges to

$$
\begin{equation*}
-\infty<\int_{D_{r} \cap A \varepsilon} \Pi_{B}(x, d z) a(y-z)<\infty \tag{8.44}
\end{equation*}
$$

as $\lambda \downarrow 0$, since $a^{\lambda}(y-z) \rightarrow a(y-z)$ uniformly for $z \in \bar{B} \cap D_{r} \cap A_{\varepsilon}^{c}$. Let $B_{r}=$ $B \cap D_{r}$, and note that $\Pi_{B_{r}}(x, d z) \geqq \Pi_{B}(x, d z)$ for $z \in D_{r}$. Since $-a^{\lambda}(y-z) \uparrow$ $-a(y-z), z \in A_{\varepsilon}$, it follows from the monotone convergence theorem that

$$
\begin{align*}
&-\int_{D_{r} \cap A_{\varepsilon}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)  \tag{8.45}\\
&= E_{x}\left(\exp \left\{-\lambda V_{B}\right\}\left[-a^{\lambda}\left(y-X_{V_{B}}\right) 1_{A_{\varepsilon}}\left(X_{V_{B}}\right) 1_{D_{r}}\left(X_{V_{B}}\right)\right)\right. \\
& \uparrow E_{x}\left[-a\left(y-X_{V_{B}}\right) 1_{A_{\varepsilon}}\left(X_{V_{B}}\right) 1_{D_{r}}\left(X_{V_{B}}\right)\right] \\
&=-\int_{D_{r} \cap A_{\varepsilon}} \Pi_{B}(x, d z) a(y-z) .
\end{align*}
$$

Since $a(y-z) \leqq 0$ for $z \in A_{\varepsilon}$, we see that

$$
\begin{equation*}
\int_{D_{r} \cap A_{\varepsilon}} \Pi_{B}(x, d z) a(y-z) \geqq \int_{D_{r} \cap A_{\varepsilon}} \Pi_{B_{r}}(x, d z) a(y-z) . \tag{8.46}
\end{equation*}
$$

By Lemma 8.3,

$$
\begin{equation*}
\int_{\bar{B}_{r}} \Pi_{B_{r}}(x, d z) a(y-z)>-\infty, \tag{8.47}
\end{equation*}
$$

and it is clear that

$$
\begin{equation*}
\infty>\int_{\bar{B}_{r} \cap\left(D_{r} \cap A_{\varepsilon}\right)^{c}} \Pi_{B_{r}}(x, d z) a(y-z)>-\infty . \tag{8.48}
\end{equation*}
$$

Hence,

$$
\begin{equation*}
\lim _{\lambda \downharpoonright 0} \int_{D_{r \cap} \cap \varepsilon} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{D_{r \cap} A_{\varepsilon}} \Pi_{B}(x, d z) a(y-z)>-\infty \tag{8.49}
\end{equation*}
$$

Finally, as in the case when $y \in(\bar{B})^{c}$, monotone convergence shows that

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \int_{D F} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{D \xi} \Pi_{B}(x, d z) a(y-z) \geqq 0 . \tag{8.50}
\end{equation*}
$$

Thus, using (8.43), we see that (8.39) holds for $y \in \bar{B}$. This completes the proof.
We may now easily establish Theorem 8.2.
Proof of Theorem 8.2. Since $B$ is nonpolar some relatively compact subset $A \subset B$ must be nonpolar. But then, by Lemma 8.3, $g_{B}(x, y) \leqq g_{A}(x, y)<\infty$ for $x \neq y$. Also

$$
\begin{align*}
L_{B}^{\lambda}(x) & =g^{\lambda}(1)\left[1-E_{x}\left(\exp \left\{-\lambda V_{B}\right\}\right)\right]  \tag{8.51}\\
& \leqq g^{\lambda}(1)\left[1-E_{x}\left(\exp \left\{-\lambda V_{A}\right\}\right)\right]=L_{A}^{\lambda}(x)
\end{align*}
$$

so

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \sup _{B}^{\lambda}(x) \leqq L_{A}(x)<\infty \tag{8.52}
\end{equation*}
$$

Using (8.5), we see that for $x \neq y$,

$$
\begin{equation*}
\lim _{\lambda \downarrow 0}\left[\int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)+L_{B}^{\lambda}(x)\right]=a(y-x)+g_{B}(x, y) \tag{8.53}
\end{equation*}
$$

has a finite limit.
Using (8.52) and Lemma 8.5, we see that

$$
\begin{equation*}
\lim _{\lambda \downharpoonright 0} \int_{\bar{B}} \Pi_{B}^{\lambda}(x, d z) a^{\lambda}(y-z)=\int_{\bar{B}} \Pi_{B}(x, d z) a(y-z), \tag{8.54}
\end{equation*}
$$

must be finite for $x \neq y$. Thus, it must be that $\lim _{\lambda \downarrow 0} L_{B}^{\lambda}(x)=L_{B}(x)$ exists and that (8.9) is satisfied. This establishes the theorem.

One of the main applications of Theorem 8.2 is to show that $g_{B}(x, y)<\infty$, $x \neq y, x, y \in B^{c}$, when $B$ is a compact nonpolar set. As we will see in the next section, $g_{B}$ restricted to $B^{c} \times B^{c}$ (for $B$ a closed set) is just the Green function of $B^{c}$. Now if one is interested only in showing that $g_{B}(x, y)<\infty$ for $x \neq y$, $x, y \in B^{c}$ and in showing that (8.9) is valid for such $x, y$ when $B$ is a nonpolar compact set then the proof of the theorem can be considerably shortened. In fact, all one needs is Lemma 8.2 and the simple parts of Lemmas 8.1 and 8.3.

A simple consequence of Theorem 8.2 is the following result.
Theorem 8.3. Let $B$ be a nonpolar Borel set. Then

$$
\begin{equation*}
\lim _{t \rightarrow \infty} \log (t) P_{x}\left(V_{B}>t\right)=2 \pi L_{B}(x) . \tag{8.55}
\end{equation*}
$$

Proof. One easily checks that

$$
\begin{equation*}
\lambda g^{\lambda}(1) \sim \frac{1}{2 \pi} \lambda \log \left(\frac{1}{\lambda}\right) \tag{8.56}
\end{equation*}
$$

The theorem follows from this and the fact that

$$
\begin{equation*}
\lim _{\lambda \downarrow 0} \lambda g^{\lambda}(1)\left[\frac{1-E_{x}\left(\exp \left\{-\lambda V_{B}\right\}\right)}{\lambda}\right] \equiv L_{B}(x) \tag{8.57}
\end{equation*}
$$

exists by well-known Tauberian theorems.
Of course Theorem 8.3 is uninteresting when $E_{x} V_{B}<\infty$ since then $L_{B}(x) \equiv 0$. When $B$ is relatively compact, however, $L_{B}(x)>0$ for all $|x|$ sufficiently large. In Section 10, we will see that $L_{B}(x)=\lim _{|y| \rightarrow \infty} g_{B}(x, y)$ whenever $B$ is relatively compact.

## 9. Green function

Now that we have Theorem 8.2 at our disposal, we can easily derive the properties of $g_{B}(x, y)$ for a nonpolar set $B$.

Theorem 9.1. Let $B$ be nonpolar. Then
(i) $0 \leqq g_{B}(x, y)<\infty$ for $x \neq y$,
(ii) $g_{B}(x, y)=g_{B}(y, x)$,
(iii) $g_{B}(x, y)+a(y-x)$ is harmonic in $y$ on $(\bar{B})^{c}$,
(iv) $g_{B}(x, y)$ is harmonic in $y$ on $(\bar{B})^{c}-\{x\}$,
(v) $g_{B}(x, y)$ is subharmonic in $y$ and upper semicontinuous in $y$ on $R^{2}-\{x\}$,
(vi) $\lim _{y \rightarrow y_{0}} g_{B}(x, y)=g_{B}\left(x, y_{0}\right)=0, y_{0} \in B^{r}$.

Proof. Part (i) is part of Theorem 8.2. Part (ii) follows from the fact that it is true for $g_{B}^{\lambda}$ and letting $\lambda \downarrow 0$. Parts (iii) to (v) follow from (8.9) and the fact that $a(y-x)$ is harmonic in $y \neq x$ and that for any finite measure $\mu$, $\int_{\bar{B}} \log |y-x| \mu(d x)$ is an upper semicontinuous and subharmonic function that is harmonic in $(\bar{B})^{c}$. Finally, (vi) follows from the upper semicontinuity and the fact that $g_{B}^{\lambda}\left(x, y_{0}\right)=0$ for all $\lambda$ and $x$.

An open set $G$ is called Greenian if there exists a function $g(x, y)$ on $G \times G$ such that $g(x, y)+a(y-x)$ is harmonic in $y$ on $G$. If $G$ is Greenian the smallest such function is called the Green function of $G$.

From our work in Section 5, we know that in dimension $n \geqq 3$ every open set is Greenian and that $g_{G^{c}}$ restricted to $G \times G$ is its Green function. Based on this, and Theorem 9.1, we would fully expect that the following holds in the planar case.

Theorem 9.2. An open set $G \subset R^{2}$ is Greenian if and only if $G^{c}$ is nonpolar. In that case $g_{G^{c}}$ restricted to $G \times G$ is the Green function of $G$.

Proof. Suppose first that $G$ is a bounded open set such that each point of $\partial G$ is regular for $G^{c}$. Then clearly $G^{c}$ is nonpolar so Theorem 9.1 shows that $g_{G^{c}}$ restricted to $G \times G$ is Greenian. To see that it is the smallest of the Greenian functions, suppose $g$ is another such function. Then by property (vi) of Theorem 9.1,

$$
\begin{equation*}
\liminf _{y \rightarrow y_{0}}\left[g(x, y)-g_{G^{c}}(x, y)\right]=\liminf _{y \rightarrow y_{0}} g(x, y) \geqq 0 \tag{9.1}
\end{equation*}
$$

Thus, by the minimum principle, $g(x, y)-g_{G^{c}}(x, y) \geqq 0$ on $G$.
Suppose now that $G$ is any open subset of $R^{2}$. By Corollary 3.1, we can find bounded open sets $G_{1} \subset \bar{G}_{1} \subset G_{2} \subset \cdots, \cup_{n} G_{n}=G$, such that each point of $\partial G_{n}$ is regular for $G_{n}^{c}$ and such that $P_{x}\left(V_{\partial G_{n}} \uparrow V_{\partial G}\right)=1$ for all $x \in G$.

By Lemma 8.4, the sequence $g_{G_{h}^{c}}$ is increasing and bounded above by $g_{G^{c}}(x, y)$. Since $g_{G_{n}}(x, y)+a(y-x)$ is harmonic in $y \in G$ for fixed $x \in G_{n}$, Harnack's theorem tells us that the limit function $g^{*}(x, y) \leqq g_{G^{c}}(x, y)$ is such that in a given component of $G$ it is either identically infinite or a harmonic function on $G-\{x\}$. Let $f \geqq 0$ be arbitrary. Then for any $x \in G$,

$$
\begin{equation*}
E_{x} \int_{0}^{V_{o G_{n}}} f\left(X_{t}\right) d t=E_{x} \int_{0}^{V_{G_{n}^{t}}} f\left(X_{t}\right) d t \uparrow E_{x} \int_{0}^{V_{G^{c}}} f\left(X_{t}\right) d t, \quad n \uparrow \infty \tag{9.2}
\end{equation*}
$$

But

$$
\begin{equation*}
E_{x} \int_{0}^{V_{G_{n}^{\epsilon}}} f\left(X_{t}\right) d t=\int_{R^{2}} g_{G_{n}^{\delta}}(x, y) f(y) d y \tag{9.3}
\end{equation*}
$$

so monotone convergence gives that

$$
\begin{equation*}
\int_{R^{2}} g_{G \hbar}(x, y) f(y) d y \uparrow \int_{R^{2}} g^{*}(x, y) f(y) d y \tag{9.4}
\end{equation*}
$$

as $n \uparrow \infty$. Thus,

$$
\begin{equation*}
\int_{R^{2}} g^{*}(x, y) f(y) d y=\int_{R^{2}} g_{G^{c}}(x, y) f(y) d y \tag{9.5}
\end{equation*}
$$

Suppose $G^{c}$ is nonpolar. Then as $g^{*}(x, y) \leqq g_{G^{c}}(x, y), g^{*}(x, y)<\infty$ for $x \neq y$, and thus $g^{*}(x, y)+a(y-x)$ is harmonic in $y$ on $G$. Since the function $f$ in (9.5) is any nonnegative function, (9.5) implies that $g^{*}(x, y)=g_{G^{c}}(x, y)$ a.e. $y$. Since $g_{G^{c}}(x, y)+a(y-x)$ is harmonic in $y$, we see that $g^{*}=g_{G^{c}}$ on $G \times G$. Suppose $g$ is any function having the required properties. Then this function restricted to $G_{n}$ also has the required properties so by what has already been proved, $g(x, y) \geqq g_{G_{\hbar}}(x, y), x, y \in G_{n}$. Thus,

$$
\begin{equation*}
g_{G^{c}}(x, y)=\lim _{n} g_{G_{\hbar}}(x, y) \leqq g(x, y) . \tag{9.6}
\end{equation*}
$$

Hence, $g_{G^{c}}$ restricted to $G \times G$ is the Green function of $G$. Finally, suppose $G^{c}$ is polar. Then $G$ cannot be Greenian. Indeed, if it were Greenian, then let $g$ be a function with the required properties. But then, as argued above

$$
g^{*}(x, y) \leqq g(x, y)<\infty, \quad x \neq y
$$

But (9.4) shows that $g^{*}(x, y)=\infty$ for a.e. $y$, a contradiction. This completes the proof.

## 10. Logarithmic potentials

Let $\mu$ be a bounded measure having compact support $k$. The function $\varphi_{\mu}(x)=$ $-\int_{k} a(y-x) \mu(d y)$ is called the potential of $\mu$. One easily verifies that $\varphi_{\mu}(x)$ is a lower semicontinuous function that is superharmonic on $R^{2}$ and harmonic on $k^{c}$.

Let $B$ be a nonpolar relatively compact set. Since $a(y-x)-a(y) \rightarrow 0$ as $|y| \rightarrow \infty$, uniformly in $x$ on compacts, it follows at once from (8.9) that, uniformly in $x$ on compacts,

$$
\begin{equation*}
\lim _{|y| \rightarrow \infty} g_{B}(x, y)=L_{B}(x) \tag{10.1}
\end{equation*}
$$

As $g_{B}(x, y)=g_{B}(y, x)$, we see that

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} g_{B}(x, y)=L_{B}(y) \tag{10.2}
\end{equation*}
$$

Let $D_{r}$ be the closed disk of center 0 and radius $r$. The hitting measure, $\Pi_{D_{r}}(x, d y)$, is just the unit mass at $x$ for $x \in D_{r}$. For $x \in D_{r}^{c}$ and $\varphi$ a continuous function on $\partial D_{r}, \Pi_{D_{r}} \varphi(x)=\Pi_{\partial D_{r}} \varphi(x)$ is the unique bounded solution to the

Dirichlet problem for $D_{r}^{c}$ with boundary function $\varphi$. As is well known from texts on complex variables or partial differential equations, the solution to this Dirichlet problem is provided by the Poisson integral. Thus,

$$
\begin{equation*}
\Pi_{D_{r}}(x, d y)=\frac{\left[|x|^{2}-r^{2}\right]}{|y-x|^{2}} \sigma_{r}(0, d y), \quad x \in D_{r}^{c} \tag{10.3}
\end{equation*}
$$

It follows from (10.3) that

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} \Pi_{D_{r}}(x, d y)=\sigma_{r}(0, d y), \tag{10.4}
\end{equation*}
$$

in the sense of strong convergence of measures.
Let $B$ be any nonpolar relatively compact set and let $D_{r}$ be a disk of center 0 and radius $r$ that contains $\bar{B}$ in its interior. Then for any bounded function $\varphi$

$$
\begin{equation*}
\Pi_{B} \varphi(x)=\int_{\partial D_{r}} \Pi_{D_{r}}(x, d z) \Pi_{B} \varphi(z), \quad x \in D_{r}^{c} \tag{10.5}
\end{equation*}
$$

Hence,

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} \Pi_{B} \varphi(x)=\int_{\partial D_{r}} \Pi_{B} \varphi(z) \sigma_{r}(0, d z) \tag{10.6}
\end{equation*}
$$

Let

$$
\begin{equation*}
\mu_{B}(d y)=\int_{\partial D_{r}} \Pi_{B}(z, d y) \sigma_{r}(0, d z) \tag{10.7}
\end{equation*}
$$

Equation (10.6) shows that $\mu_{B}(d y)=\lim _{|x| \rightarrow \infty} \Pi_{B}(x, d y)$ in the sense of strong convergence of measures. We have thus proved the following important result.

Theorem 10.1. Let $B$ be a nonpolar relatively compact set. Then $g_{B}(x, y) \rightarrow$ $L_{B}(y)$ as $|x| \rightarrow \infty$, and $\Pi_{B}(x, d y) \rightarrow \mu_{B}(d y)$ as $|x| \rightarrow \infty$ in the sense of strong convergence.

The measure $\mu_{B}$ has the obvious probabilistic significance as the hitting probability of $B$ starting from infinity. We will show that $\mu_{B}$ should be considered in potential theoretic terms as the equilibrium measure of $B$.

Theorem 10.2. Let $B$ be a nonpolar relatively compact set. Then

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty}\left[L_{B}(x)-a(x)\right]=k(B) \tag{10.8}
\end{equation*}
$$

exists and is finite. Moreover,

$$
\begin{equation*}
\varphi_{\mu_{B}}(x)=k(B)-L_{B}(x) \tag{10.9}
\end{equation*}
$$

Proof. Suppose $y \notin \bar{B}$. Then $a(y-z)$ is bounded for $z \in \bar{B}$, and thus by Theorem 10.2,

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} \int_{\bar{B}} \Pi_{B}(x, d z) a(y-z)=\int_{\bar{B}} \mu_{B}(d z) a(y-z) \tag{10.10}
\end{equation*}
$$

exists and is finite. Using (8.9), we see that

$$
\begin{align*}
& a(y-x)-a(x)-\int_{\bar{B}} \Pi_{B}(x, d z) a(y-z)+g_{B}(x, y)  \tag{10.11}\\
&=L_{B}(x)-a(x) .
\end{align*}
$$

Using (10.10) and (10.2), we see that the left side converges to $\varphi_{\mu_{B}}(y)+L_{B}(y)$ as $|x| \rightarrow \infty$. Hence, the right side must have a finite limit. This establishes (10.8) and (10.9) for $y \in(\bar{B})^{c}$. Suppose $y \in \bar{B}$. Then (8.9) shows

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} \int_{\vec{B}} \Pi_{B}(x, d z) a(y-z)=L_{B}(y)-k(B) . \tag{10.12}
\end{equation*}
$$

The function $g_{B}(\xi, y)+L_{B}(\xi)-a(y-\xi)$ is clearly bounded in $\xi$ if $|\xi|>r$ for some sufficiently large $r$. But then by (8.9),

$$
\begin{equation*}
\int_{\bar{B}} \Pi_{B}(\xi, d z) a(y-z)=g_{B}(\xi, y)+L_{B}(\xi)-a(y-\xi) \tag{10.13}
\end{equation*}
$$

is bounded in $\xi$ for $|\xi|>r$. Now for any closed disk $D_{r}$ of center 0 and radius $r$ containing $\bar{B}$ in its interior,

$$
\begin{equation*}
\int_{\bar{B}} \Pi_{B}(x, d z) a(y-z)=\int_{\partial D_{r}} \Pi_{D_{r}}(x, d \xi) \int_{\bar{B}} \Pi_{B}(\xi, d z) a(y-z), \quad x \notin \bar{D} . \tag{10.14}
\end{equation*}
$$

Thus, by Theorem 10.1 and equation (10.7),

$$
\begin{align*}
\lim _{|x| \rightarrow \infty} \int_{\bar{B}} \Pi_{B}(x, d z) a(y-z) & =\int_{\partial D_{r}} \sigma_{r}(0, d \xi) \int_{\bar{B}} \Pi_{B}(\xi, d z) a(y-z)  \tag{10.15}\\
& =\int_{\bar{B}} \mu_{B}(d z) a(y-z) .
\end{align*}
$$

This establishes the theorem for $y \in \bar{B}$ and thereby completes the proof.
Definition 10.1. Let $B$ be a nonpolar relatively compact set. The measure $\mu_{B}$ in Theorem 10.1 is called the equilibrium measure of $B$. The constant $k(B)$ in Theorem 10.2 is called the Robin's constant of $B$ and the potential of $\mu_{B}$ is called the equilibrium potential of $B$.

For a relatively compact polar set, we define $k(B)=+\infty$. Theorem 10.3 given below will show that this is the natural definition of $k(B)$ for a polar set.

Proposition 10.1. Let $A$ and $B$ be two relatively compact sets such that $A \subset B$. Then $k(A) \geqq k(B)$.

Proof. If $B$ is polar then $A$ must also be polar. In this case $k(A)=k(B)=$ $+\infty$. Suppose $B$ is nonpolar. If $A$ is polar, then $k(A)=+\infty$ and $k(B)<\infty$ so the proposition is valid. Suppose $A$ is also nonpolar. Then $L_{A} \geqq L_{B}$ (since $L_{A}^{\lambda} \geqq L_{B}^{\lambda}$ ) and ( 10.8 ) shows that $k(A) \geqq k(B)$.

Theorem 10.3. Let $B$ be a compact set. Then

$$
\begin{equation*}
k(B)=\sup \{k(U): U \text { open, } U \supset B \text { and } \bar{U} \text { compact }\} . \tag{10.16}
\end{equation*}
$$

On the other hand, if $U$ is an open relatively compact set, then

$$
\begin{equation*}
k(U)=\inf \{k(A): A \text { compact, } A \subset U\} \tag{10.17}
\end{equation*}
$$

Proof. Suppose $B$ is compact. Let $B_{n}, n \geqq 1$ be relatively compact open sets such that $B_{1} \supset \bar{B}_{2} \supset B_{2} \supset \cdots, \cap_{n} B_{n}=\cap_{n} \bar{B}_{n}=B$. Then as was shown in Proposition 3.2, $P_{x}\left(V_{B_{n}} \uparrow V_{B}\right)=1$ for all $x \in B^{c} \cup B^{r}$. Suppose that $B$ is polar. Let $f \geqq 0$ be continuous with compact support and have integral 1 and set $A f(z)=\int_{R^{2}} a(y-z) f(y) d y$. Then for any $x \in B^{c}$,

$$
\begin{equation*}
\int_{R_{n}} g_{B_{n}}(x, y) f(y) d y=E_{x} \int_{0}^{V_{B_{n}}} f\left(X_{t}\right) d t \uparrow \infty, \quad n \rightarrow \infty \tag{10.18}
\end{equation*}
$$

Now $A f(x)$ is a continuous function, and thus for all $n$

$$
\begin{equation*}
\left|\int_{\bar{B}_{n}} \Pi_{B_{n}}(x, d z) A f(z)\right| \leqq \sup _{z \in \overline{B_{1}}}|A f(z)|=M<\infty \tag{10.19}
\end{equation*}
$$

Using (8.9), we see that

$$
\begin{align*}
\left|-\int_{R^{n}} g_{B_{n}}(x, y) f(y) d y+L_{B_{n}}(x)\right| & =\left|A f(x)-\int_{\bar{B}_{n}} \Pi_{B_{n}}(x, d z) A f(z)\right|  \tag{10.20}\\
& \leqq|A f(x)|+M<\infty
\end{align*}
$$

Thus, using (10.18), we see that $L_{B_{n}}(x) \uparrow \infty$ for each $x \in B^{c}$. By (10.9),

$$
\begin{equation*}
\int_{R^{2}} \varphi_{\mu_{B_{n}}}(x) f(x) d x+\int_{R^{2}} L_{B_{n}}(x) f(x) d x=k\left(B_{n}\right) \tag{10.21}
\end{equation*}
$$

But

$$
\begin{equation*}
\left|\int_{R^{2}} \varphi_{\mu_{B_{n}}}(x) f(x) d x\right|=\left|\int_{R^{2}} A f(x) \mu_{B_{n}}(d x)\right| \leqq M<\infty \tag{10.22}
\end{equation*}
$$

and $\int_{R^{2}} L_{B_{n}}(x) f(x) d x \uparrow \infty$. Thus, $k\left(B_{n}\right) \uparrow \infty$. This establishes (10.16) when $B$ is polar.

Suppose now that $B$ is nonpolar. Let $f$ and $A f$ be as before. Then for $x \in B^{c} \cup B^{r}$,

$$
\begin{align*}
\int_{R^{n}} g_{B_{n}}(x, y) f(y) d y & =E_{x} \int_{0}^{V_{B_{n}}} f\left(X_{t}\right) d t \uparrow E_{x} \int_{0}^{V_{B}} f\left(X_{t}\right) d t  \tag{10.23}\\
& =\int_{R^{2}} g_{B}(x, y) f(y) d y .
\end{align*}
$$

Also as $P_{x}\left(X_{V_{B_{n}}} \in \bar{B}_{1}\right)=1$ for all $n$.

$$
\begin{align*}
\lim _{n \rightarrow \infty} \int_{\bar{B}_{n}} \Pi_{B_{n}}(x, d z) A f(z) & =\lim _{n \rightarrow \infty} E_{x} A f\left(X_{V_{B_{n}}}\right)  \tag{10.24}\\
& =E_{x} A f\left(X_{V_{B}}\right)=\int_{\bar{B}} \Pi_{B}(x, d z) A f(z) .
\end{align*}
$$

Hence by (8.9), for $x \in B^{c} \cup B^{r}$,
(10.25)

$$
\begin{aligned}
\lim _{n \rightarrow \infty} & L_{B_{n}}(x) \\
& =\lim _{n \rightarrow \infty}\left[A f(x)-\int_{\bar{B}_{n}} \Pi_{B_{n}}(x, d z) A f(z)+\int_{R^{2}} g_{B_{n}}(x, y) f(y) d y\right] \\
& =A f(x)-\int_{\bar{B}} \Pi_{B}(x, d z) A f(z)+\int_{R^{2}} g_{B}(x, y) f(y) d y=L_{B}(x) .
\end{aligned}
$$

Since $\left(B^{r}\right)^{c} \cap B$ has measure 0 , we see that

$$
\begin{equation*}
\int_{R^{2}} L_{B_{n}}(x) f(x) d x \uparrow \int_{R^{2}} L_{B}(x) f(x) d x \tag{10.26}
\end{equation*}
$$

Now if $D$ is a disk of center 0 containing $\bar{B}$ in its interior,

$$
\begin{align*}
\lim _{n \rightarrow \infty} \int_{R^{2}} \mu_{B_{n}}(d x) A f(x) & =\lim _{n \rightarrow \infty} \int_{\partial D} \sigma(0, d \xi) E_{\xi} A f\left(X_{V_{B_{n}}}\right)  \tag{10.27}\\
& =\int_{\partial D} \sigma(0, d \xi) E_{\xi} A f\left(X_{V_{B}}\right) \\
& =\int_{R^{2}} \mu_{B}(d x) A f(x)
\end{align*}
$$

Hence, using this fact, (10.26), and (10.21), we see that $k\left(B_{n}\right) \uparrow k(B)$. This establishes (10.16). To prove (10.17), note that we can find compacts $A_{n} \subset U$ such that $A_{1} \subset A_{2} \subset \cdots, \cup_{n} A_{n}=U$. But then $P_{x}\left(V_{A_{n}} \downarrow V_{U}\right)=1$ for all $x$. The remainder of the proof of (10.17) is similar to the proof of (10.16) for $B$ a nonpolar set. We omit these details.

Let $A$ and $B$ be two Borel sets. Then,

$$
\begin{align*}
P_{x}\left(V_{A \cap B} \leqq t\right) & \leqq P_{x}\left(V_{A} \leqq t, V_{B} \leqq t\right)  \tag{10.28}\\
& \leqq P_{x}\left(V_{A} \leqq t\right)+P_{x}\left(V_{B} \leqq t\right)-P_{x}\left(V_{A \cup B} \leqq t\right) .
\end{align*}
$$

Thus, $P_{x}\left(V_{A \cap B}>t\right) \geqq P_{x}\left(V_{A}>t\right)+P_{x}\left(V_{B}>t\right)-P_{x}\left(V_{A \cup B}>t\right)$. It follows from this and (8.6) that

$$
\begin{equation*}
L_{A \cap B}^{\lambda}(x) \geqq L_{A}^{\lambda}(x)+L_{B}^{\lambda}(x)-L_{A \cup B}^{\lambda}(x) . \tag{10.29}
\end{equation*}
$$

Letting $\lambda \downarrow 0$, we see that whenever $A$ and $B$ are nonpolar

$$
\begin{equation*}
L_{A \cap B}(x)+L_{A \cup B}(x) \geqq L_{A}(x)+L_{B}(x) . \tag{10.30}
\end{equation*}
$$

If we take $L_{k}(x) \equiv \infty$ whenever $k$ is polar, then (10.30) is valid for all sets.
Using (10.8), we see that for relatively compact sets

$$
\begin{equation*}
-k(A \cup B)+[-k(A \cap B)] \leqq-k(A)+(-k(B)) \tag{10.31}
\end{equation*}
$$

Also as $L_{A}^{\lambda}(x) \geqq L_{B}^{\lambda}(x)$ for $A \subset B$, we see that $L_{A}(x) \geqq L_{B}(x)$, so for $A \subset B$,

$$
\begin{equation*}
-k(A) \leqq-k(B) \tag{10.32}
\end{equation*}
$$

Define $k^{*}(B)=-k(B)$ if $B$ is compact and define $k^{*}(U)=\sup \left\{k^{*}(B): B \subset U\right.$, $B$ compact $\}$ if $U$ is open. Then (10.16), (10.17), (10.31), and (10.32) show that $k^{*}(\cdot)$ is a Choquet capacity on the compact sets. The extension theorem of Choquet then implies that for any Borel set $B$,

$$
\begin{equation*}
\sup \left\{k^{*}(A): A \subset B, A \text { compact }\right\}=\inf \left\{k^{*}(U): U \supset B, U \text { open }\right\} \tag{10.33}
\end{equation*}
$$

and that the common value $k^{*}(B)$ is such that $k^{*}(A) \leqq k^{*}(B), A \subset B$ and

$$
\begin{equation*}
k^{*}(A \cup B)+k^{*}(A \cap B) \leqq k^{*}(A)+k^{*}(B) \tag{10.34}
\end{equation*}
$$

Let $B$ be any relatively compact set. Then for any compact set $A \subset B$ and any open set $U \supset B,-k(A) \leqq-k(B) \leqq-k(U)$. Thus,

$$
\begin{equation*}
k^{*}(B) \leqq-k(B) \leqq k^{*}(B) \tag{10.35}
\end{equation*}
$$

Equation (10.35) tells us that whenever $B$ is relatively compact, $k(B)=$ $-k^{*}(B)$. For a general Borel set, define $k(B)$ to be $-k^{*}(B)$ and define the capacity $C(B)$ of $B$ to be $e^{-k(B)}$. Our discussion above has shown the following result.

Theorem 10.4. Let $B$ be any relatively compact set. Then

$$
\begin{equation*}
\inf \{k(A): A \subset B, A \text { compact }\}=k(B)=\sup \{k(U): U \text { open, } U \supset B\} . \tag{10.36}
\end{equation*}
$$

Corollary 10.1. A Borel set $B$ has capacity 0 if and only if it is polar.
Proof. The set $B$ has capacity 0 if and only if $k(B)=\infty$. We have already shown that for a relatively compact set $B$ this is the case if and only if $B$ is polar. If $B$ is not relatively compact and $k(B)=\infty$, then (10.36) shows that $k(A)=\infty$ for every compact subset of $B$, and thus (again by (10.36)) $k(D)=\infty$ for every relatively compact subset $D$ of $B$. But then every relatively compact subset of $B$ is polar, and as $B$ is a countable union of relatively compact sets, $B$ must be polar. Conversely, if $B$ is polar, every compact subset of $B$ is polar, so for any compact subset $A \subset B, k(A)=\infty$. Thus, (10.36) shows that $k(B)=\infty$. This establishes the corollary.

If $B$ is a relatively compact set in dimension $n \geqq 3$, then as was shown in Section 6, $B$ has capacity 0 (that is, $B$ is polar) if and only if the only finite measure $\mu$ having support on $B$ with a bounded potential is the 0 measure. For potentials in the plane we have the following analog.

Theorem 10.5. A relatively compact set B is polar if and only if $\sup _{x} \varphi_{\mu}(x)=$ $+\infty$ for every nonzero finite measure $\mu$ having support on $B$.

Proof. Let $A$ be a relatively compact open set containing $\bar{B}$. Let $a_{N}(x)=$ $a(x)$ if $a(x) \geqq-N$ and let $a_{N}(x)=-N$ if $a(x)<-N, N>0$. Then clearly

$$
\begin{equation*}
-\int_{R^{2}} \mu(d x) \int_{R^{2}} a_{N}(y-x) \mu_{A}(d y)=-\int_{R^{2}} \mu_{A}(d y) \int_{R^{2}} a_{N}(y-x) \mu(d x) \tag{10.37}
\end{equation*}
$$

By monotone convergence as $N \uparrow \infty$, we see that

$$
\begin{equation*}
\int_{\bar{B}} \mu(d x) \varphi_{\mu_{A}}(x)=\int_{\bar{A}} \mu_{A}(d x) \varphi_{\mu}(x) . \tag{10.38}
\end{equation*}
$$

Since $\bar{B} \subset A$ equation (10.9) shows that the left side is $\mu(\bar{B}) k(A)$. Thus, $\mu(\bar{B}) k(A) \leqq \sup _{x} \varphi_{\mu}(x)$, and hence by Theorem 10.4, $\mu(\bar{B}) k(B) \leqq \sup _{x} \varphi_{\mu}(x)$. If $B$ is polar then $k(B)=+\infty$. On the other hand, if $\sup _{x} \varphi_{\mu}(x)=+\infty$ for all nonzero $\mu$ supported on $\bar{B}$, then $B$ must be polar. For, if $B$ were nonpolar then the equilibrium measure $\mu_{B}$ of $B$ would be a probability measure on $B$ whose potential would be $\leqq k(B)<+\infty$ everywhere, a contradiction. This establishes the theorem.

Remark. By using the maximum principle for potentials $\varphi_{\mu}$ (see, for example, Hille, Analytic Function Theory, Vol. II), we replace $\sup _{x}$ by $\sup _{x \in \bar{B}}$ in Theorem 10.5.


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